



# $O_2$ in the 565-700 nm region



## FT-IR measurements of the $O_2$ - $O_2$ collision-induced absorptions (cia) at high pressures

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## □ Key optics components

(1) A broad-band spectrometer: JPL Bruker IFS-125HR

(2) A (heatable) high-pressure cell (HHPC) developed: Rated to 150 bars at 520 K

(3) A broad-band super-luminous light source: A Laser Driven Light Source.

**All three components have been configured and employed for this work!**

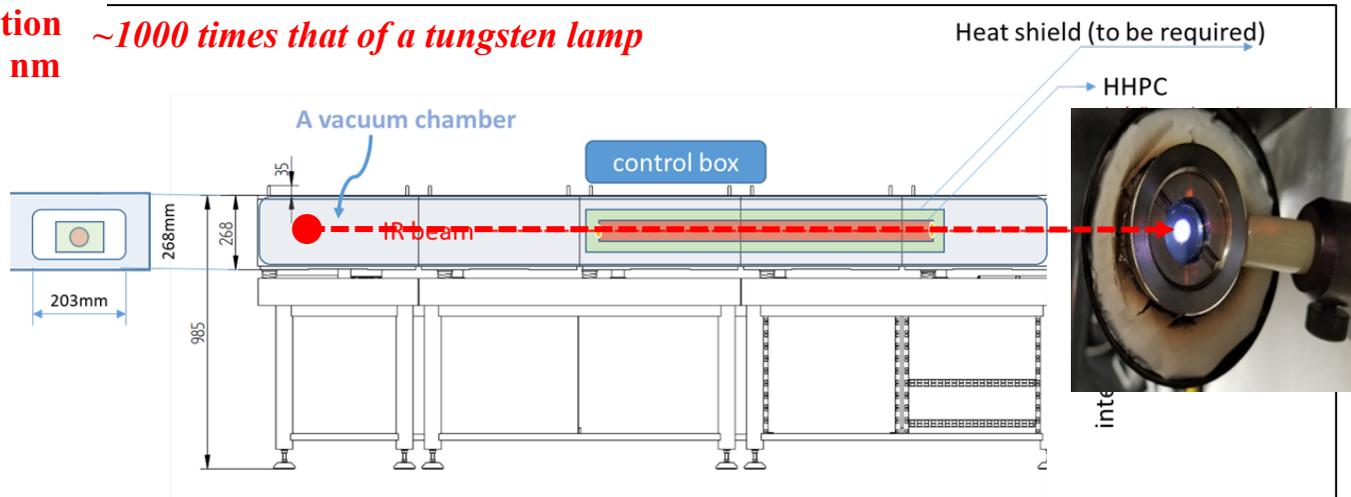
## □ HHPC inside an extension chamber

**Broadband collimation  
in the 200 – 2400 nm**

*~1000 times that of a tungsten lamp*



EQ-99FC LDLS with Fiber Collimator





# Experimental conditions

## - Number of molecules in the beam path



### • Gas samples

- O<sub>2</sub>, 99.999%, normal sample [ $\varepsilon < 0.1\%$ ]
- Dry air, 21(1)% O<sub>2</sub> [ $\varepsilon \sim 4.8\%$ ]

### • Path 1: HHPC

- Cell path = 0.999(1) m
- Prs and Temp [ $\delta=0.1$  bar and 0.3K, respectively]
- Pressure ranges, 3 – 131 bars
- Departure from ideal gas law [up to 5.5%]

Pure O <sub>2</sub> [bars]	correction factors	Dry air	correction factors
116.4	1.055	131.2	1.043
83.0	1.045	130.9	1.043
57.3	1.034	101.8	1.033

### • Other paths (FT-IR & Ext. Chambers, Source box)

### ♣ Virial Eq. of State for real gases

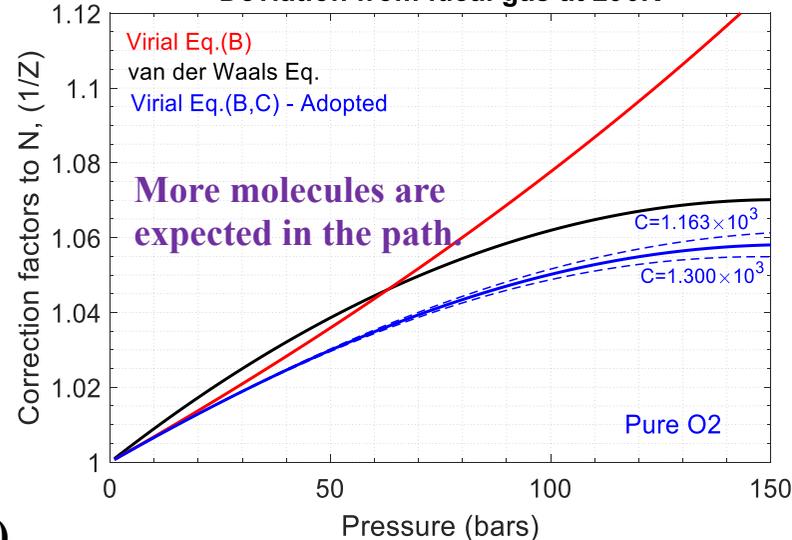
The compressibility factor,  $Z$ , for a gas mixture is described by

$$Z = \frac{PV_m}{RT} = 1 + \frac{B_{mix}}{V_m} + \frac{C_{mix}}{V_m^2} + \dots$$

For a mixture,

$$B_{mix} = x_1^2 B_{11} + 2x_1 x_2 B_{12} + x_2^2 B_{22}$$

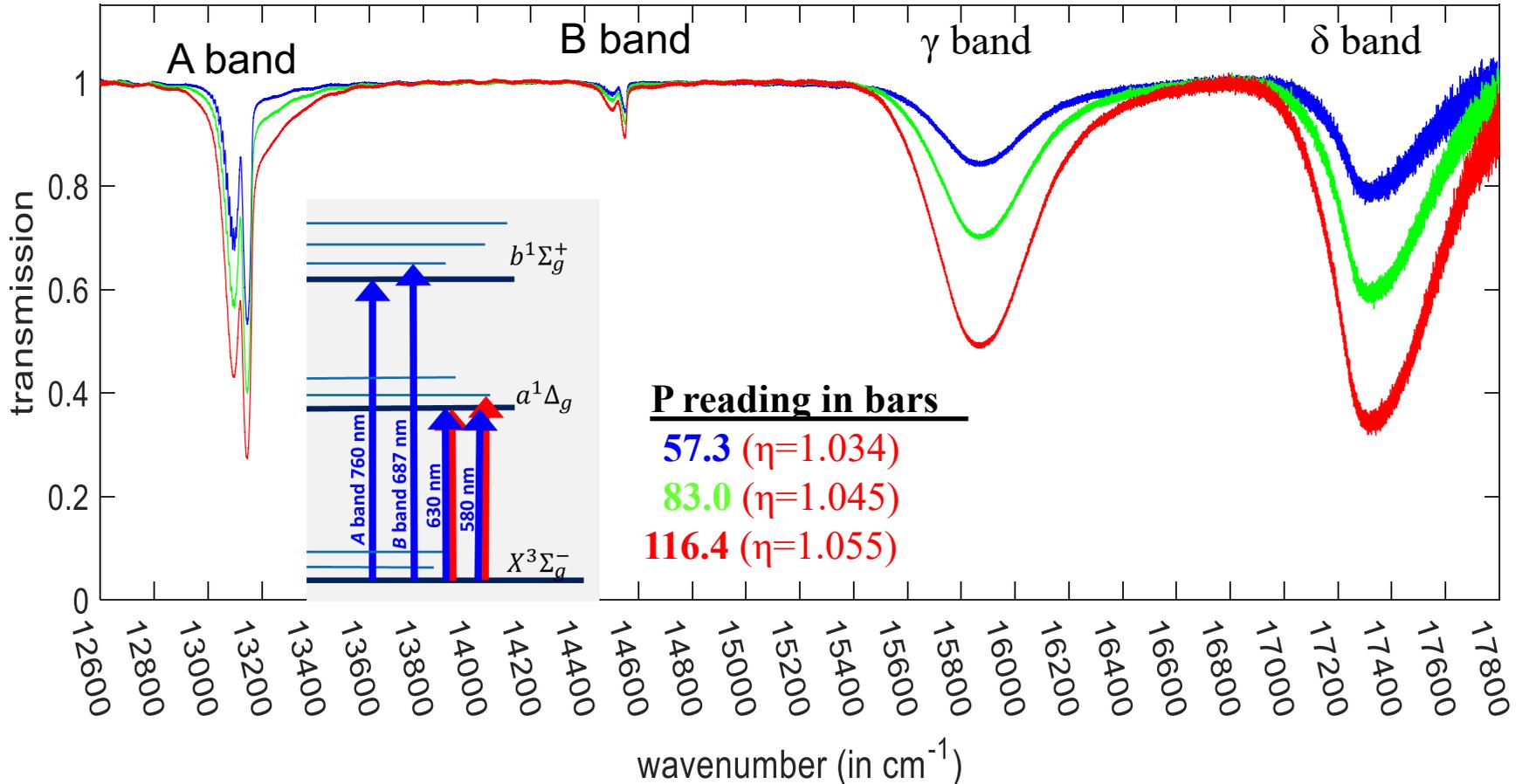
Deviation from ideal gas at 296K



Note: The 2<sup>nd</sup>, 3<sup>rd</sup> Virial constants,  $B$  and  $C$ , are adopted from the CRC Handbook.

# Spectra obtained at room Temp.

Pure O<sub>2</sub> at high pressures





# How to derive cia for O<sub>2</sub>(B, γ, δ)



## ➤ What is measured:

### Total extinction

RS = Rayleigh scattering (→ Ratioed out)

MD = magnetic dipole moments (Simulated)

- $\alpha_{\text{tot}} = \alpha_{\text{RS}} + \alpha_{\text{MD}} + \alpha_{\text{cia}}$
- $\alpha_{\text{MD}} = \alpha_{\text{O}_2\_lines} (+ \alpha_{\text{O}_2\_LM})$

## ➤ $\alpha_{\text{MD}}$ : Model & Assumptions:

### (1) Widths and shifts at high pressures

→ Assumed proportional to Prs.

### (2) Profile: SD-Voigt+LM (HITRAN)

Speed-dependence parameters  
from Domyslawska et al.(2016)

### (3) $\alpha_{\text{LM}}$ : Line mixing parameters

Rosenkranz parameters

Not available for O<sub>2</sub>(B) band  
*so we calculated them.*

## ➤ Line mixing parameter calculations

- $K$  = inelastic state-to-state collisional rotational transition (or transfer) rates;
- Upward transition modelled by EPG Law

$$K_{j \leftarrow k}^{EPG} = a \left| \frac{\Delta E_{jk}}{B} \right|^{-b} \exp\left(-c \left| \frac{\Delta E_{jk}}{k_B T} \right| \right)$$

- Downward transition by Detailed balance

$$K_{k \leftarrow j}^{EPG} = \frac{\rho_k}{\rho_j} K_{j \leftarrow k}^{EPG}$$

- $W$ , Real part of Relaxation Matrix

$$W_{jk} = -\frac{\epsilon}{2} [K_{j \leftarrow k}^{EPG} + K_{j' \leftarrow k'}^{EPG}]$$

- Model validation by Observed widths

$$\gamma_l(\text{obs}) = \ll lkk' | \text{Re}[W] | l'j' \gg = W_{kk'}$$

- Finally, first-order line mixing coefficient,

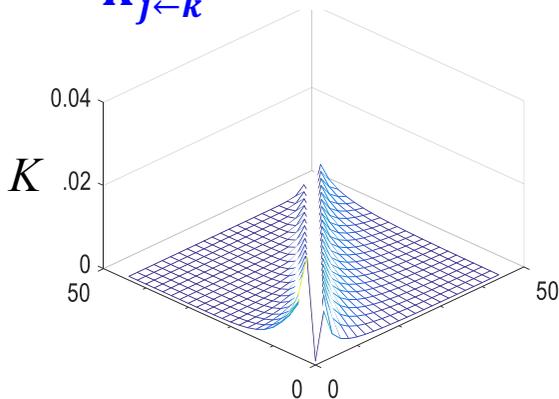
**Rosenkranz parameter,  $Y$**

$$Y_k = 2 \sum_{j \neq k} \frac{d_j}{d_k} \frac{W_{jk}}{v_j - v_k}$$

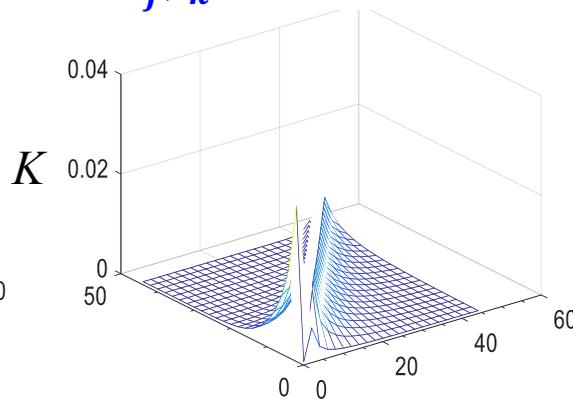
# Observed line widths $\rightarrow [W] \rightarrow$ Line mixing coeffs.

Treat PP, PQ, RQ, RR branches, separately.

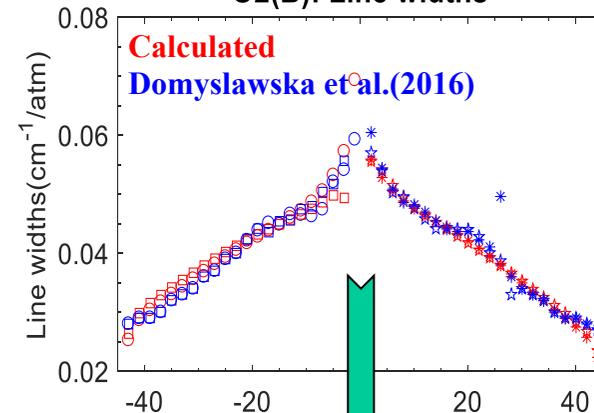
$K_{j \leftarrow k}^{EPG}$  Lower vib. state



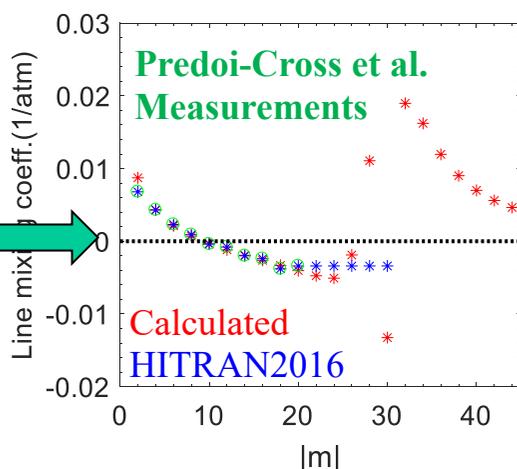
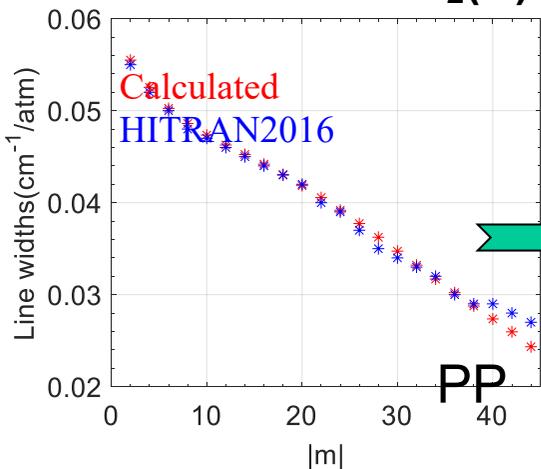
$K_{j \leftarrow k}^{EPG}$  Upper vib. state



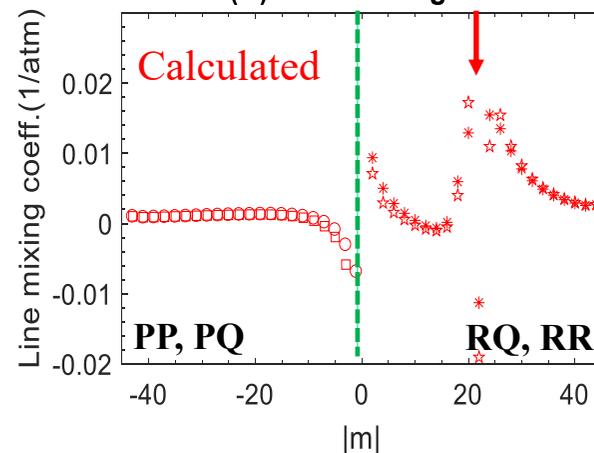
O2(B): Line widths



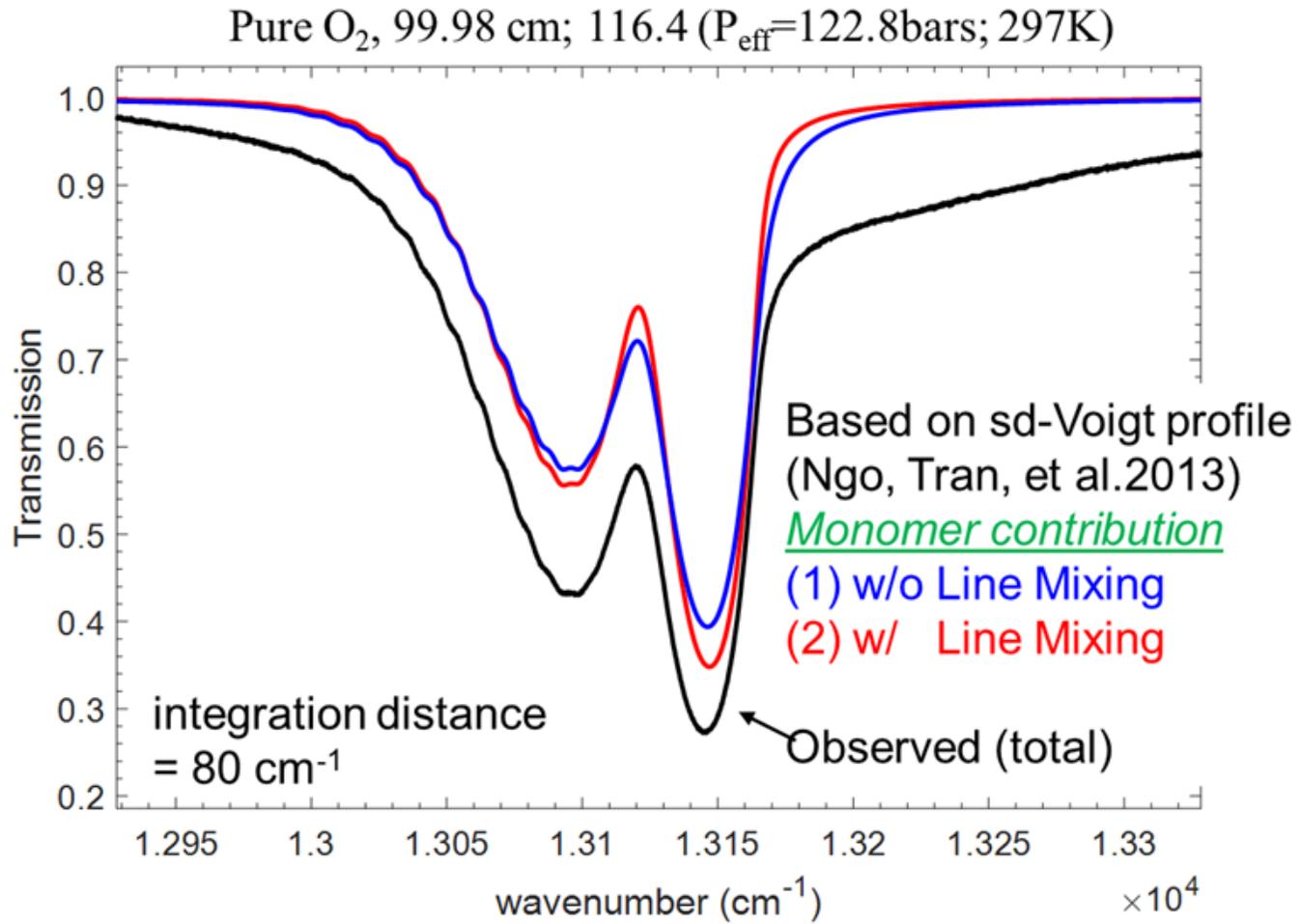
O2(A) RR-branch



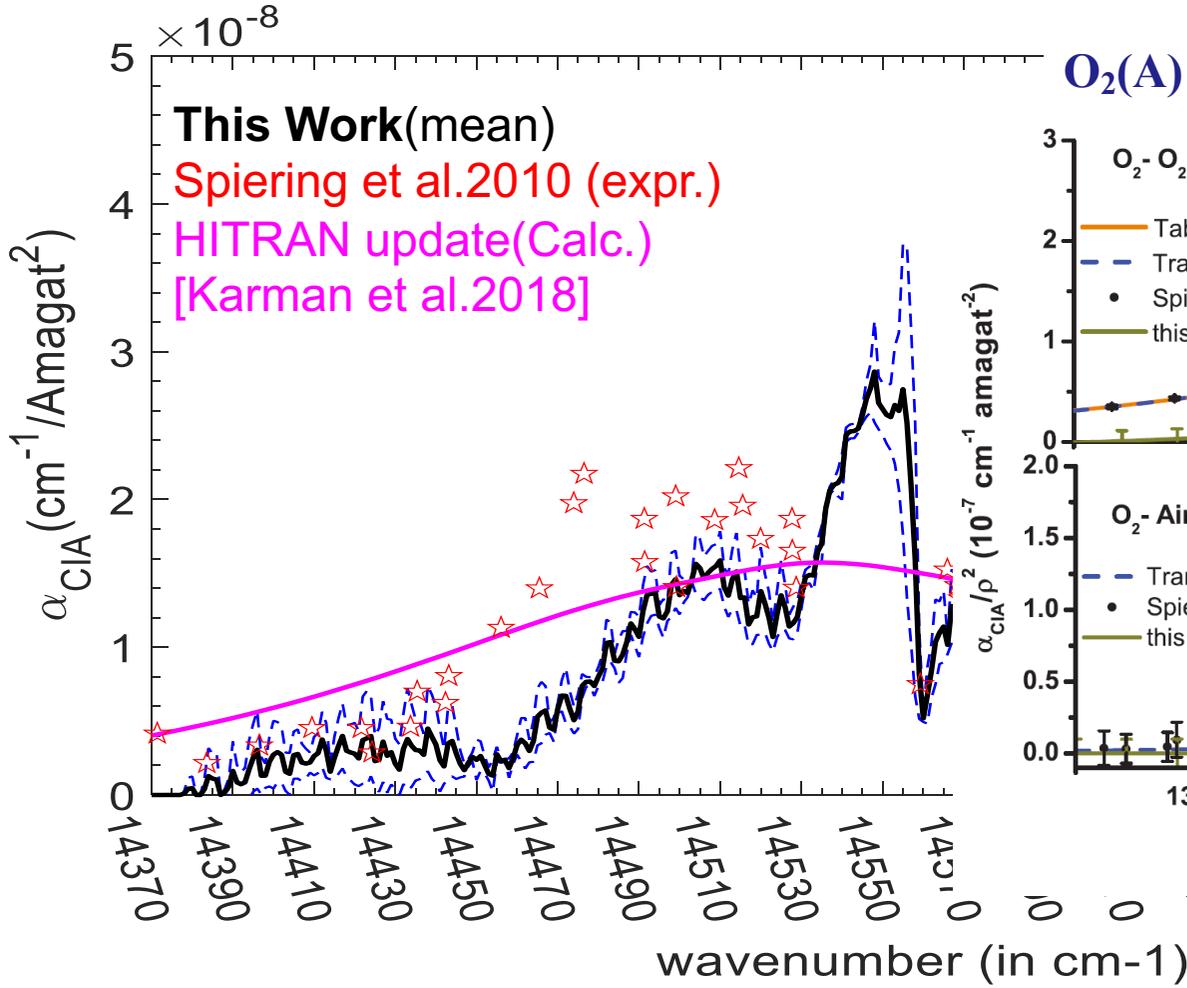
O2(B): Line Mixing coeff.



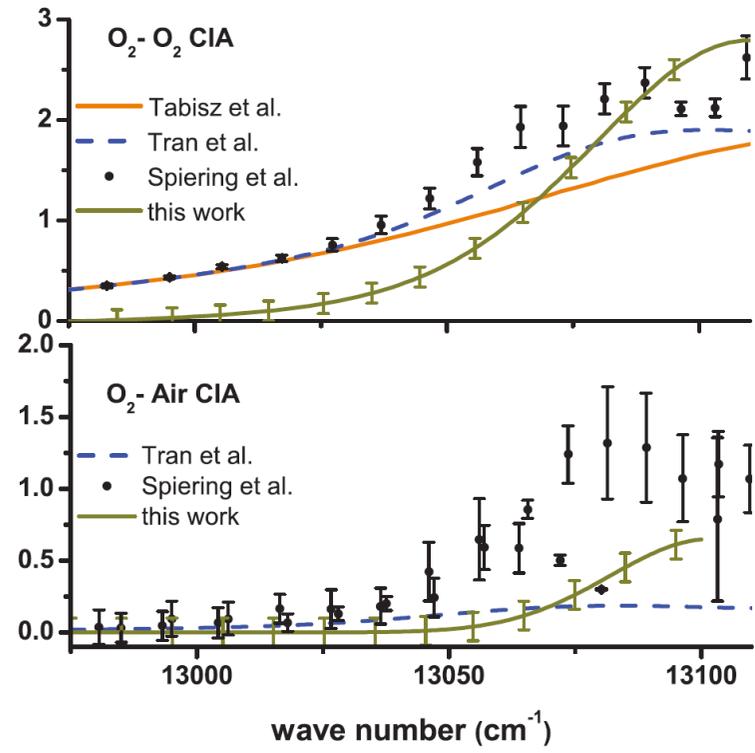
# Impact of Line mixing – a sample case

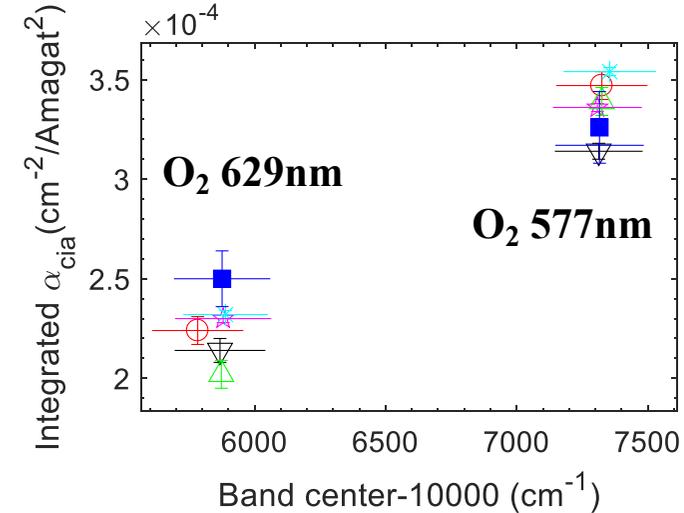
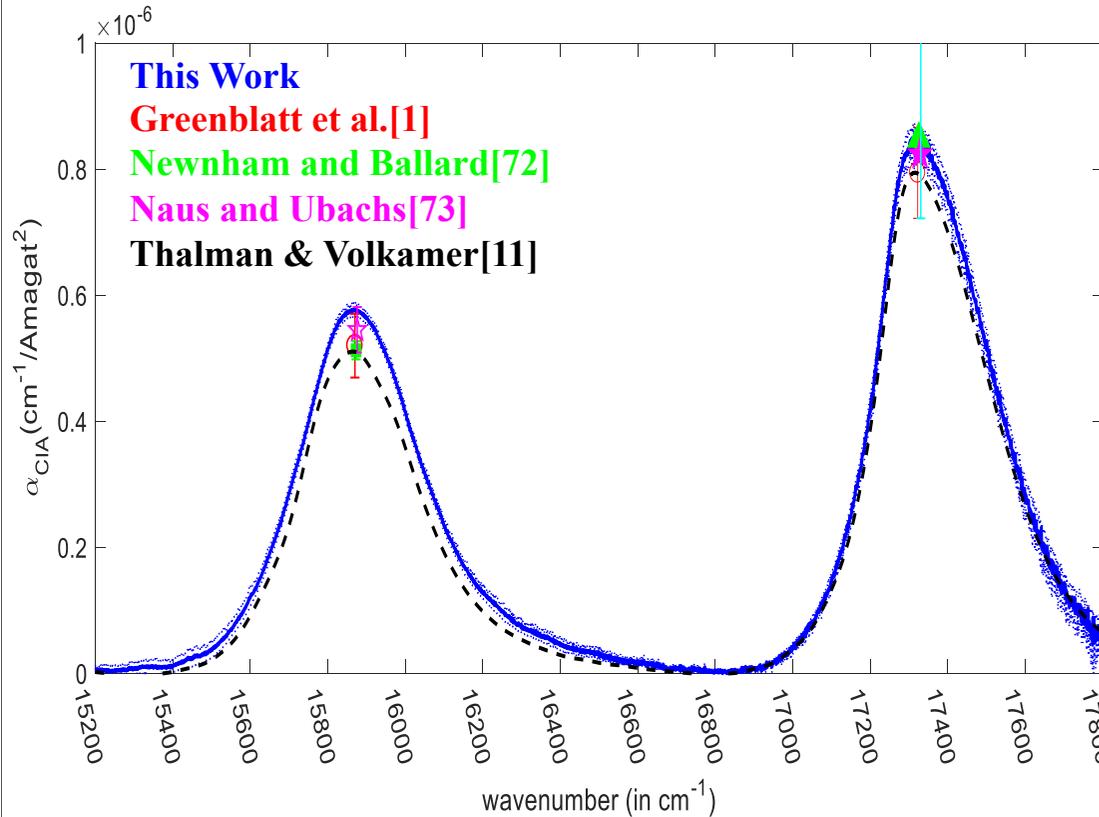


# CIA in the O<sub>2</sub>(B) band based on sd-Voigt + 1st order line mixing (Rosenkranz)



## O<sub>2</sub>(A) band cia [Long et al. 2012]



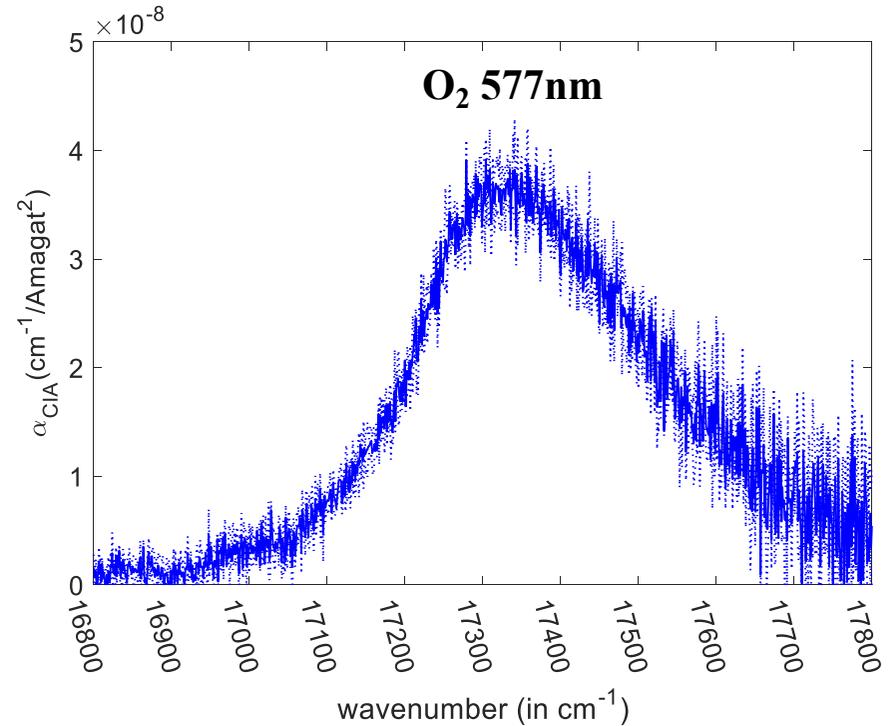
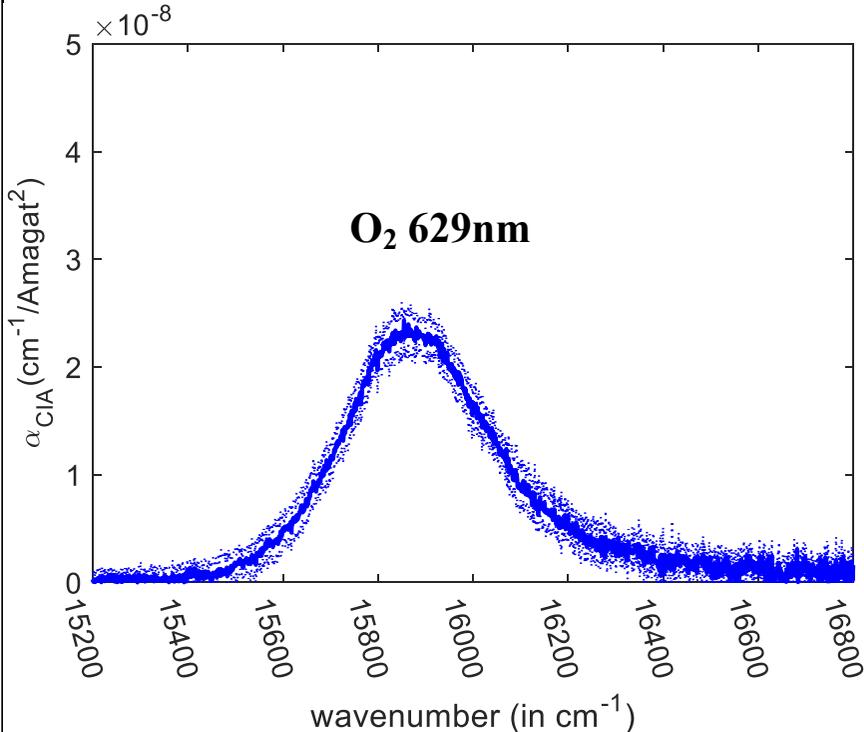


- Comparison:
- Peak frequency
  - FWHM,
  - Integrated CIA

- (1) Agreed within combined uncertainties
- (2) This work shows higher than others.
- (3) Band shapes are different.

- Comparison:
- Spread, 629 and 577 nm bands
  - Band center ~ 106 & 43 cm<sup>-1</sup>
  - FWHM ~ 47 & 18 cm<sup>-1</sup>
  - Integrated CIA ~ 8% & 4%

- Weaker than that of O<sub>2</sub>-O<sub>2</sub> cia by a factor of ~30
- No measurements available for O<sub>2</sub>-N<sub>2</sub> cia



Regions	T (K)	a <sup>1</sup> Δ <sub>g</sub> (v = 0) + a <sup>1</sup> Δ <sub>g</sub> (v = 0)	a <sup>1</sup> Δ <sub>g</sub> (v = 1) + a <sup>1</sup> Δ <sub>g</sub> (v = 0)
Work			
band centers		629 nm	577 nm
Present work	296K	1.03(±0.21)×10 <sup>-5</sup>	1.48(±0.20)×10 <sup>-5</sup>



# Summary and Acknowledgements



## ☐ Measured O<sub>2</sub>-O<sub>2</sub> CIA in the O<sub>2</sub>-B, γ, and δ band regions

- Simulate the resonance contribution with Line Mixing taken into account
- O<sub>2</sub>(B): Integrated CIA =  $0.32(6) \times 10^{-5} \text{ cm}^{-2}/\text{Amg}^2$   
cf. Spiering et al.(2011):  $S_{\text{cia}} = 0.44(4) \times 10^{-5} \text{ cm}^{-2}/\text{Amg}^2$
- Ratio O<sub>2</sub>:  $\text{cia}(A)/\text{cia}(B) = 16.7$  cf. A-band/B-band = 16.9  
cf. Tran et al. (2006) Ratio = 10.7 ; Spiering et al.(2011), Ratio = 14.9
- **Integrated O<sub>2</sub>-O<sub>2</sub> cia in the O<sub>2</sub>- γ and δ band regions:**  
 $25.0(1.4) \times 10^{-5}$  and  $32.6(1.8) \times 10^{-5}$ , respectively.
- **Integrated O<sub>2</sub>-N<sub>2</sub> cia in the O<sub>2</sub>- γ and δ band regions:**  
 $1.03(0.21) \times 10^{-5}$  and  $1.48(0.20) \times 10^{-5}$ , respectively.

☐ Full line mixing element calculations for O<sub>2</sub>(A) – on the way.

## Acknowledgements

**OCO-2, IMAP, NASA-ACLS program, JPL R&TD for Air Quality Res. for support;  
Drs. M. Devi, D.C Benner, A. Predoi-Cross, Ha Tran for helpful discussion.**