

## Fabrication of High- $T_c$ Hot-Electron Bolometric Mixers for Terahertz Applications

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**Abstract** - Superconducting hot-electron bolometers (HEB) represent a promising candidate for heterodyne mixing at frequencies exceeding 1 THz. Nb HEB mixers offer performance competitive with tunnel junctions without the frequency limit imposed by the superconducting energy gap. Although the performance of  $\text{YBa}_2\text{Cu}_3\text{O}_{7-x}$  HEB mixers is not projected to be superior to that of Nb devices, which operate at low temperatures, they introduce the possibility of sensitive, low power heterodyne detectors operating at temperatures approaching 90 K for applications requiring portability and closed-cycle refrigeration. We report on the fabrication and characterization, both DC and RF, of high- $T_c$  mixers based on ultra-thin (<20 nm)  $\text{YBa}_2\text{Cu}_3\text{O}_{7-x}$  films patterned to micrometer dimensions and incorporated into 2.5 THz planar mixer circuits.

### 1. INTRODUCTION

Heterodyne receivers based on superconductor-insulator-superconductor (SIS) junctions are used in submillimeter astronomical and atmospheric studies at frequencies up to 1 THz. At higher frequencies, Schottky diode mixers or direct detection techniques are generally used. Superconducting hot-electron bolometers (HEB), consisting of a thin film intimately coupled to a cooled substrate, have been proposed for use in mixing applications in the terahertz (THz) range [1,2]. Superconducting HEB mixers require orders of magnitude less local oscillator (LO) power than semiconductor mixers, allowing for solid state LOs rather than the large lasers required of Schottky diode mixers. Nb mixers operating at 2-4 K have shown excellent performance at 530 GHz [3], and recently demonstrated low noise at 1.2 THz [4] and 2.5 THz [5]. While some initial experiments on device physics and heterodyne performance [6,7] have been carried out on high-temperature superconductor (HTS) bolometers, no optimized devices have been reported to date.

The lower LO power requirements of superconducting HEB mixers makes them particularly attractive for remote

sensing systems with constraints on the availability of power or weight, such as on balloon or stratospheric-based platforms. HEB mixers made from HTS operating in the range of 60-85 K are particularly attractive for space based applications since they can be cooled with existing space qualified closed-cycle refrigerators.

### II. REQUIREMENTS

The operating principles of a 2.5 THz HTS HEB mixer have been described in detail elsewhere [8]. HEB mixer operation depends on breaking heating the electrons with incoming radiation, resulting in a non-equilibrium energy distribution. Cooling occurs via electron-phonon interactions whereby the hot electrons give their energy to phonons which escape into the substrate. Alternatively, cooling can occur by diffusion of the hot electrons out of the device and into the normal-metal electrical contacts [2]. Due to the short electron mean-free-path in HTS materials, the electron-phonon cooling mechanism dominates in HTS HEB mixers. In space applications, system-level power restrictions constrain mixers to operate at low LO power levels, with major implications for device size.

As a result of these considerations, the major design requirements [8] for HTS HEB's are described below. The first four requirements pertain to the substrate. While a large number of substrates can individually meet these conditions, it is important for optimal mixer performance to meet all simultaneously.

(1) The substrate must have a high thermal conductivity and be compatible with epitaxial  $\text{YBa}_2\text{Cu}_3\text{O}_{7-x}$  (YBCO) growth. This requirement can be met by a number of substrates compatible with high-quality YBCO films.  $\text{MgO}$ ,  $1-x\text{AlO}_3$ ,  $\text{Al}_2\text{O}_3$  and  $\text{YAlO}_3$  have thermal conductivity's at 90 K of 3.4, 0.35, 6.4 and  $0.2-0.4 \text{ W K}^{-1} \text{ cm}^{-1}$ , respectively. For comparison, yttrium stabilized zirconia (YSZ) has a thermal conductivity at 90 K of only  $0.015 \text{ W K}^{-1} \text{ cm}^{-1}$  [8].

(2) The Kapitza boundary resistance ( $R_b$ ) between the HTS film and the substrate should be as small as possible. This second requirement can be met by several of the aforementioned substrates. The values of  $R_b$  inferred from measurements of the phonon escape time, between YBCO and  $\text{MgO}$ ,  $1-x\text{AlO}_3$  and  $\text{Al}_2\text{O}_3$  at 90 K are 5, 10, and  $11 \times 10^{-4} \text{ K cm}^2 \text{ W}^{-1}$  respectively [9-13].

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(3) Substrates need to have a small loss-tangent at both 2.5 THz and at the IF frequency. The fourth requirement can also be met by a number of the aforementioned substrates. MgO, LaAlO<sub>3</sub>, Al<sub>2</sub>O<sub>3</sub> and YAlO<sub>3</sub> have loss tangents of 7, 5, 8, and 10,  $\times 10^{-6}$ , respectively at 90 K and  $\sim 10$  GHz. For comparison, yttrium stabilized zirconia (YSZ) has a loss tangent of  $400 \times 10^{-4}$  at 90 K at  $\sim 10$  GHz [8].

(4) Substrates need to have a convenient dielectric constant at both 2.5 THz and at the IF frequency. This fifth requirement can also be met by a number of HTS compatible substrate materials. The dielectric constants of MgO, Y<sub>2</sub>O<sub>3</sub>, Si-on-Al<sub>2</sub>O<sub>3</sub> and YAlO<sub>3</sub>, measured in a Fourier transform spectrometer at JPL, are 10.0, 12.9, 9.9 and 21.2, respectively, at 77 K and 2.5 THz.

(5) The HTS mixer film volume must be small enough to allow device operation at microwatt (10) power levels. This third requirement dictates that the mixer be constructed from ultra-thin (10-20 nm) films patterned to micron or sub-micron dimensions [8].

In this work,  $\sim 20$  nm YBCO thick films on YAlO<sub>3</sub> substrates were used to fabricate HEB mixers.

### III. GROWTH AND PATTERNING

Growth of the superconductor and gold (Au) contact layers are performed completely *in situ* without exposure of interfaces to the ambient environment. The devices are grown on  $2.50 \mu\text{m}$  thick,  $1 \times 1 \text{cm}^2$  (001) YAlO<sub>3</sub> substrates polished on both sides. The nominal growth process is: The substrates are mounted on Haines alloy plates using Ag paint. These are transferred into the HTS deposition system via a load-lock. The substrates are buffered using a 20 nm PrBa<sub>2</sub>Cu<sub>3</sub>O<sub>7- $\delta$</sub>  (PBCO) layer deposited by pulsed laser deposition (PLD) at 790°C, 400 mTorr of O<sub>2</sub>, at a fluence of 1.6 J/cm<sup>2</sup> at  $\lambda = 248 \text{nm}$ . Substrate heating is radiative, monitored by a thermocouple that is cross checked by an optical pyrometer prior to film growth. The PBCO layer is followed by a 20 nm YBCO layer deposited at 810°C, 200 mTorr of O<sub>2</sub>, and 1.6 J/cm<sup>2</sup>. The deposited bilayer is cooled *in situ* at 40°C/minute in a 500-650 Torr O<sub>2</sub> atmosphere from the growth temperature down to room temperature. Then 100 nm of Au is deposited *in situ* by DC magnetron sputtering in a 1 mTorr Ar atmosphere (Fig. 1A). Typical transition temperature (T<sub>c</sub>) for these trilayers, as determined by AC susceptibility, is 83-86 K with a transition width of less than 2 K.

After the trilayer growth process, the substrate is mechanically removed from the Haines alloy plate. Photoresist (AZ5214) is spun onto the blank trilayers at 3000 rpm to a thickness of 1.5  $\mu\text{m}$  and soft-baked at 95 °C for 2 minutes. The sintered Ag paste residue remaining from substrate mounting on the Haines plate is scraped off of the substrate, followed by swabbing with 1000A HNO<sub>3</sub> and rinsing in water. The photoresist is then removed with acetone and the

devices are rinsed in 100% ethanol and blow-dried with dry N<sub>2</sub>.

The initial patterning of the PBCO/YBCO/Au trilayer into the antenna, RF filter, IF/DC contacts and bolometer microbridge is performed using optical contact lithography. Photoresist (AZ5206) is spun onto the unpatterned trilayers at 5000 rpm to a thickness of 0.5  $\mu\text{m}$  and soft-baked at 95°C for 2 minutes. The resist is exposed through a chrome contact mask for approximately 15 seconds at a 350 nm UV flux of 10 mW/cm<sup>2</sup>, and then developed for approximately 20 seconds in AZ developer diluted 1:1 with water. The resist is soft-baked again at 95°C for 1 minute and ashed for 30 seconds in a 40 mTorr oxygen plasma in a Semi Group 1000TP reactive ion etch (RIE) tool at 120 watts, with a DC substrate self-bias of -320 volts. The minimum feature size with good definition for this process is 1  $\mu\text{m}$ .

Next, the devices are placed into the load-lock of the deposition system in which an ion mill is located. The etching process uses normally incident 500 eV Ar<sup>+</sup> ions at 1 mA/cm<sup>2</sup> for 5 minutes (Fig 1B). The pressure is  $2.0 \times 10^{-4}$  Torr. The substrates are not cooled, however the temperature remains below 100 °C during the etching process (Fig. 1c).

After milling, the devices are transferred from the load-lock directly into the deposition system where 60 nm of YSZ is deposited at room temperature by PLD (Fig. 1D).

The photoresist is then removed by ultrasonically cleaning the devices in acetone for 1-2 minutes. The devices are rinsed in 100% ethanol and blow dried with dry N<sub>2</sub>. This YSZ deposition and lift-off process leaves the side-walls of the device coated with a protective layer of YSZ. (Fig. 1E).

At this point in the process, another layer of photoresist (AZ5206) is spun on the device to a thickness of 0.5  $\mu\text{m}$ , soft-baked, exposed, and developed using the procedure described above. The mask for this step opens a small window in the resist, exposing the 1  $\mu\text{m}$ -wide bolometer bridge, which is still covered with 100 nm of Au. The device is then placed in the RIE system for the following procedure (Fig 1F): (1) Oxygen ashing for 5 minutes in a 200 mTorr oxygen plasma at 60 watts with approximately -80 self bias. (2) Etching for 30 minutes in a 200 mTorr 1:10 O<sub>2</sub>: CCl<sub>2</sub>F<sub>2</sub> plasma at 60 watts with approximately -20 volts self bias. The approximate Au removal rate is 2.5 nm/minute. We have found that over etching does not damage c-axis-oriented YBCO. (3) Oxygen ashing for 2 minutes in a 200 mTorr oxygen plasma at 30 watts with approximately -40 volts self bias. The resulting structure is illustrated in Fig. 1G. It should be noted that without the YSZ side-wall coating covering the *a-b* plane edges of the YBCO layer, lines as wide as 50  $\mu\text{m}$  are no longer superconducting after the Au RIE process, presumably due to chlorine being driven into the film along the *a-b* planes. With the YSZ side-wall coating, we have successfully used this 3-step Au removal process on 20 nm thick YBCO lines as narrow as 400 nm and maintained T<sub>c</sub> above 80 K.

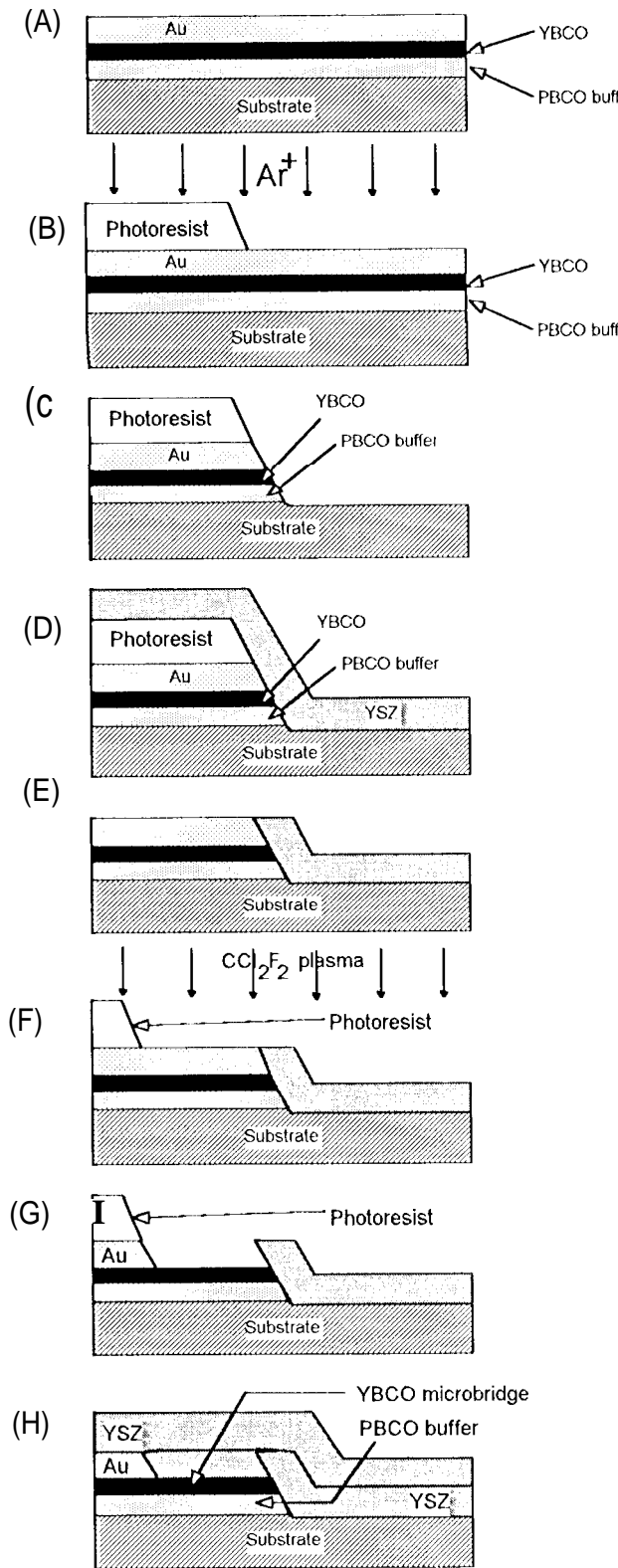


Fig. 1. Cross section of the mixer at various stages of the fabrication process. The steps are described in the text.

Next the devices are placed in the PLD system load-lock, pumped down and immediately transferred into the

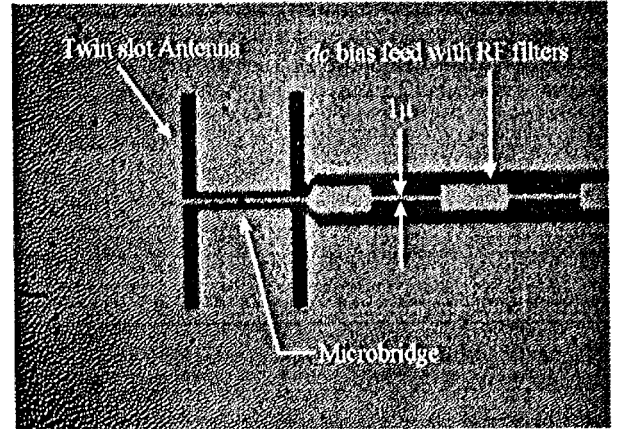


Fig. 2 - Finished HTS bolometric mixer showing the microbridge, twin slot antenna, and RF filter structure for feeding the DC bias in and extracting the intermediate frequency (IF) out.

deposition chamber', where 100 nm of YSZ is deposited by PLD, filling in the area where the Au was just removed. The total time from removal from the RIF to pumping down in the load lock is always less than 10 minutes, and typically less than 5 minutes. The photoresist is then removed by ultrasonically cleaning the devices in acetone for 1-2 minutes. The devices are rinsed in 100% ethanol and blow dried with dry N<sub>2</sub>.

The devices are next placed back into the deposition system where 100 nm of YSZ is deposited by PLD onto the entire substrate at room temperature (Fig. 11 1).

Photoresist (AZ5214) is next spun onto the 1 x 1 cm<sup>2</sup> substrates (now containing 19 chips) at 3000 rpm to a thickness of 1.5 μm and soft-baked at 95 °C for 2 minutes. The substrates are mounted on a dicing saw and cut into individual 1 x 1.5 mm<sup>2</sup> HTS chips. The photoresist is then removed with acetone and the devices are rinsed in 100% ethanol and blow dried with dry N<sub>2</sub>. A photograph of a finished device is shown in Fig. 2.

#### IV. ELECTRICAL TESTS

For DC device tests, individual dice are mounted in 28-pin Keocera ceramic chip packages which plug into the bottom of a cryogenic dipping probe. Four wires are ultrasonically bonded to the YSZ-covered Au contacts. With proper settings, the wire bonds make contact through the nominally 100 nm thick YSZ over-layer. The wire bond connections are chosen in order to allow 4-terminal device measurements, eliminating resistance contributions from the probe and instrumentation wiring.

Resistance versus temperature measurements are taken using a computer controlled system. A Keithley 220 DC current source connected to two leads applies ±1 μA and the voltage response of the device under test (DUT) is measured for both current polarities using an HP 3457A multimeter. The difference is used to eliminate contributions from ther-

## CONCLUSION

We have designed and fabricated superconducting hot-electron bolometers based on a previously developed model [4]. The devices utilize ultra-thin YBCO ( $\leq 20$  nm) films patterned into  $1 \mu\text{m}$  by  $1 \mu\text{m}$  microbridges and passivated with YSZ. These bridges maintain  $T_c$  on the order of 80 K and  $J_c$ 's  $> 1 \times 10^6 \text{ A/cm}^2$  at 77 K. We have demonstrated that the devices were successfully coupled to a 2.5 THz I.O source when operated at 77 K.

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## REFERENCES

- [1] E.M. Gershenzon, G.N. Gol'tsman, I.G. Goghidze, Y.I. Gusev, A.I. Elant'ev, H. S. Karasik and A.D. Semenov, "Millimeter and submillimeter range mixer based on electronic heating of superconducting films in the resistive state", *Sov. J. Supercond.* 3, 1582 (1990)
- [2] D.E. Prober, "Superconducting terahertz mixer using a transition-edge microbolometer", *Appl. Phys. Lett.* 62, 2119 (1993)
- [3] A. Skalare, W.R. McGrath, B. Bumble, H.G. LeDuc, J.J. Burke, A.A. Verheijen, R.J. Schoelkopf, and D.E. Prober, "Large bandwidth and low-noise in a diffusion-cooled hot-electron bolometer mixer", *Appl. Phys. Lett.* 68, 1558 (1996)
- [4] A. Skalare, W.R. McGrath, B. Bumble, and H.G. LeDuc, This conference proceedings.
- [5] B. Karasik, M. Gaidis, W.R. McGrath, B. Bumble, and H.G. LeDuc, This conference proceedings.
- [6] V.A. Trifonov, B.S. Karasik, M.A. Zorin, G.N. Gol'tsman, E.M. Gershenzon, M. Lindgren, M. Danerud, D.M. Winkler, "9.6mm wavelength mixing in a patterned  $\text{YBa}_2\text{Cu}_3\text{O}_{7.8}$  thin film", *Appl. Phys. Lett.* 68, 1418 (1996)
- [7] Yu. P. Gousev, A.D. Semenov, E.V. Pechen, A. V. Varlashkin, R.S. Nebosich, and K.F. Renk, "Coupling of terahertz radiation to a high- $T_c$  superconducting hot electron bolometer mixer", *Appl. Phys. Lett.* 69, 1 (1996)
- [8] B. Karasik, W.R. McGrath, M. Gaidis, M.J. Burns, K. Delin, A. Kleinsasser, R. Vasquez, "Modeling and optimization of a high- $T_c$  hot-electron superconducting mixer for terahertz applications", *Proceedings of the Seventh International Symposium on Space Terahertz Technology*, Charlottesville, VA, March 12-14, 1996, pp. 565-583
- [9] G.L. Carr, M. Quijada, D.B. Tanner, C.J. Hishumugi, G.P. Williams, S. Estemand, B. Dutta, F. DeRosa, A. Inam, V. Venkatesan, and X.X. Xi, "Fast bolometric response by high- $T_c$  detectors measured with subnanosecond synchrotron radiation", *Appl. Phys. Lett.* 57, 2725 (1990)
- [10] N. Bluzer, "Temporal relaxation of nonequilibrium in Y-Ba-Cu-O measured from transient photoimpedance response" *Phys. Rev. B* 44, 10222 (1991)
- [11] C.D. Marshall, J.M. Fishman, R.C. Dorfman, C.B. Eom, and M.D. Fayer, "Thermal diffusion, interfacial thermal barrier, and ultrasonic propagation in  $\text{YBa}_2\text{Cu}_3\text{O}_{7.5}$  thin films: surface-selective transient-grating experiments", *Phys. Rev. B* 45, 10009 (1992)
- [12] A.V. Seigeev, A.D. Semenov, P. Kouminov, V. Trifonov, I.G. Goghidze, B.S. Karasik, G.N. Gol'tsman, and E.M. Gershenzon, "Transparency of a  $\text{YBa}_2\text{Cu}_3\text{O}_{7.8}$  film/substrate interface for thermal phonons measured by means of voltage response to radiation", *Phys. Rev. B* 49, 9091 (1994)
- [13] M. Danerud, D. Winkler, M. Lindgren, M. Zorin, V. Trifonov, B.S. Karasik, G.N. Gol'tsman, and E.M. Gershenzon, "Nonequilibrium and bolometric photoresponse in patterned  $\text{YBa}_2\text{Cu}_3\text{O}_{7.8}$  thin films", *J. Appl. Phys.* 76, 1902 (1994)
- [14] H. S. Karasik, M.C. Gaidis, W.R. McGrath, M.J. Burns, and A.W. Kleinsasser, To be published.

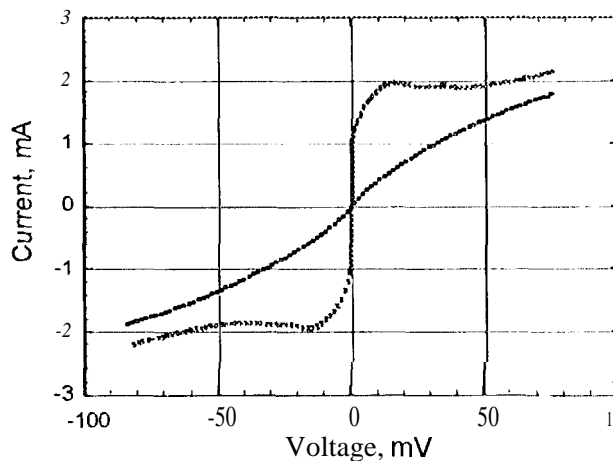
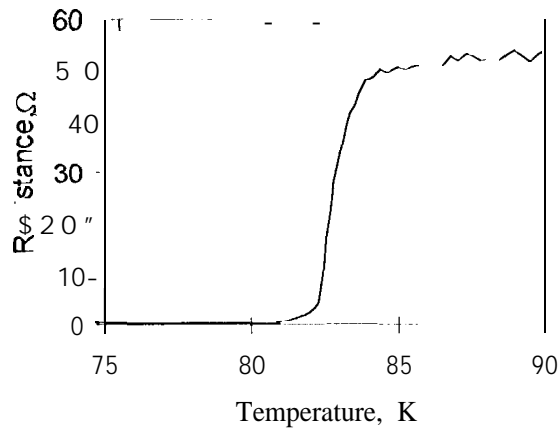


Fig. 3 - Resistance versus temperature, and I-V @ 77 K with and without the application of 2.5 THz local oscillator power, for a  $1 \times 1 \times 0.02 \mu\text{m}^3$  microbridge IEB device fabricated using the process described above.

really induced voltages in the probe and instrumentation wiring.

After processing, the microbridge  $T_c$  determined from the  $R$  vs.  $T$  data is generally 2-3K lower than the  $T_c$  of the initial trilayer film measured by AC susceptibility. The transition widths are also slightly broader. Fig. 3 shows  $R$  vs.  $T$  for a  $1 \times 1 \times 0.02 \mu\text{m}^3$  microbridge in the mixer circuit show in Fig. 2.

The samples are RF tested by mounting in an aluminum block and placed in an optical cryostat. [14] The 2.5 THz I.O consists of a methanol far-infrared laser, pumped by a  $\lambda = 9.6 \mu\text{m}$   $\text{CO}_2$  laser.  $I$ - $V$  curves at 77 K, with and without applied I.O power (estimated to be about a  $\mu\text{W}$ ), are also shown. With the application of local oscillator power, the critical current of the microbridge at 77 K can be almost entirely suppressed from its initial value of approximately  $5 \times 10^6 \text{ A/cm}^2$ .