

**MEASURING THE THICKNESS AND ELASTIC PROPERTIES
OF ELECTROACTIVE THIN-FILM POLYMERS USING
PLATEWAVE DISPERSION DATA**

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INTRODUCTION

- Electroactive thin-film polymers are offering unique capabilities for sensors and actuators.
- Muscle mechanisms and micro-electro mechanical systems (MEMS) are emerging technologies that increasingly needing such thin film polymers.
- The obtain actuation/sensing capability piezoelectric, electrostrictive or electrostatic effects are employed.
- The films are used in the thickness range of tens to hundreds of microns and the strain can be linearly or quadratically proportional to the electric field.
- In addition to the thickness change. the film vibrates as a plate structure.
- Measuring the thickness and its change under activation of an electric field and distinguishing between the thickness value and the film vibration amplitude is a complex and costly problem.
- Most methods, such as interferometry, eddy-current and capacitance, are measuring the location of the top surface of the film assuming that the rear surface stays stationary.

TECHNOLOGY NEED

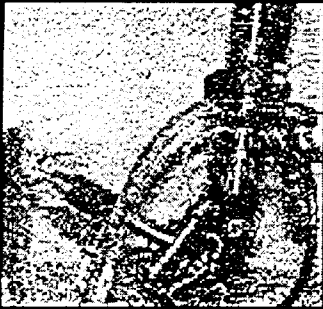
- The determination of the thickness of thin films simultaneously with the position of their surfaces (that move as a result of vibration) cannot practically be made with conventional methods.
- Electrostrictive polymers encounter internal polymer restructuring leading to potential change of the elastic properties.
- Knowledge of the elastic constants and their change is critical to the understanding of the electrostriction phenomena and the ability to distinguish it from electrostatic activation.

ULTRASONICS AS A CHARACTERIZATION TOOL

- . Ultrasonic pulse-echo offers an ideal tool for simultaneous determination of the location of the top surface, i.e., vibration amplitude, and the film thickness.
- . Obtaining an acceptable resolution for 50 to 100- μm thick films requires frequencies in the range of 50-MHz and above, which is beyond the capability of conventional ultrasonic systems.
- . Plate wave measurements allow to determine the thickness of thin films using much lower frequencies and to obtain significantly higher resolution.
- . Using dispersion curve measurements one can also determine the elastic constants of the film.

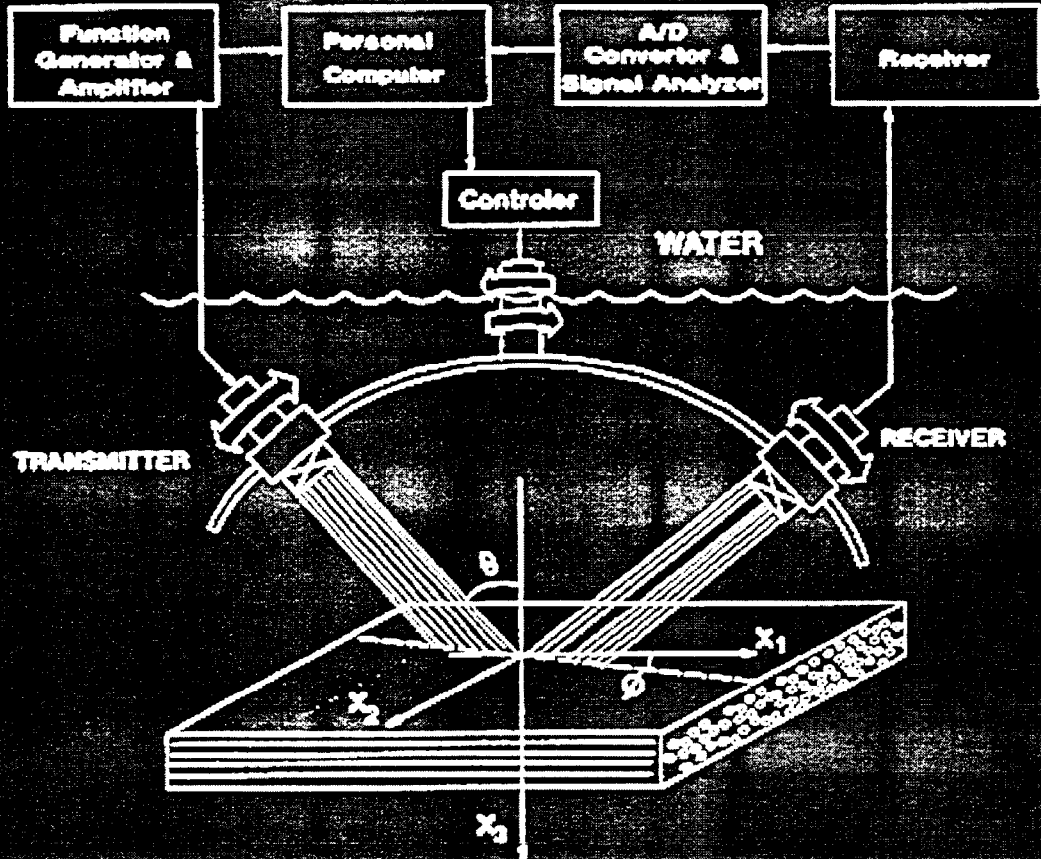
EXPERIMENTAL SETUP

- A pair of transmitter/receiver transducers is used in a pitch-catch arrangement.
- The specimen film is immersed in coupling medium and the launched wave is monitored by the receiver.
- The amplitude spectra of the reflected wave as a function of frequency is used to determine the dispersion curve of the leaky guided waves generated by the specimen.
- The dispersion curves are strongly affected by the film thickness and its elastic constants.



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LEAKY LAMB WAVE EXPERIMENT



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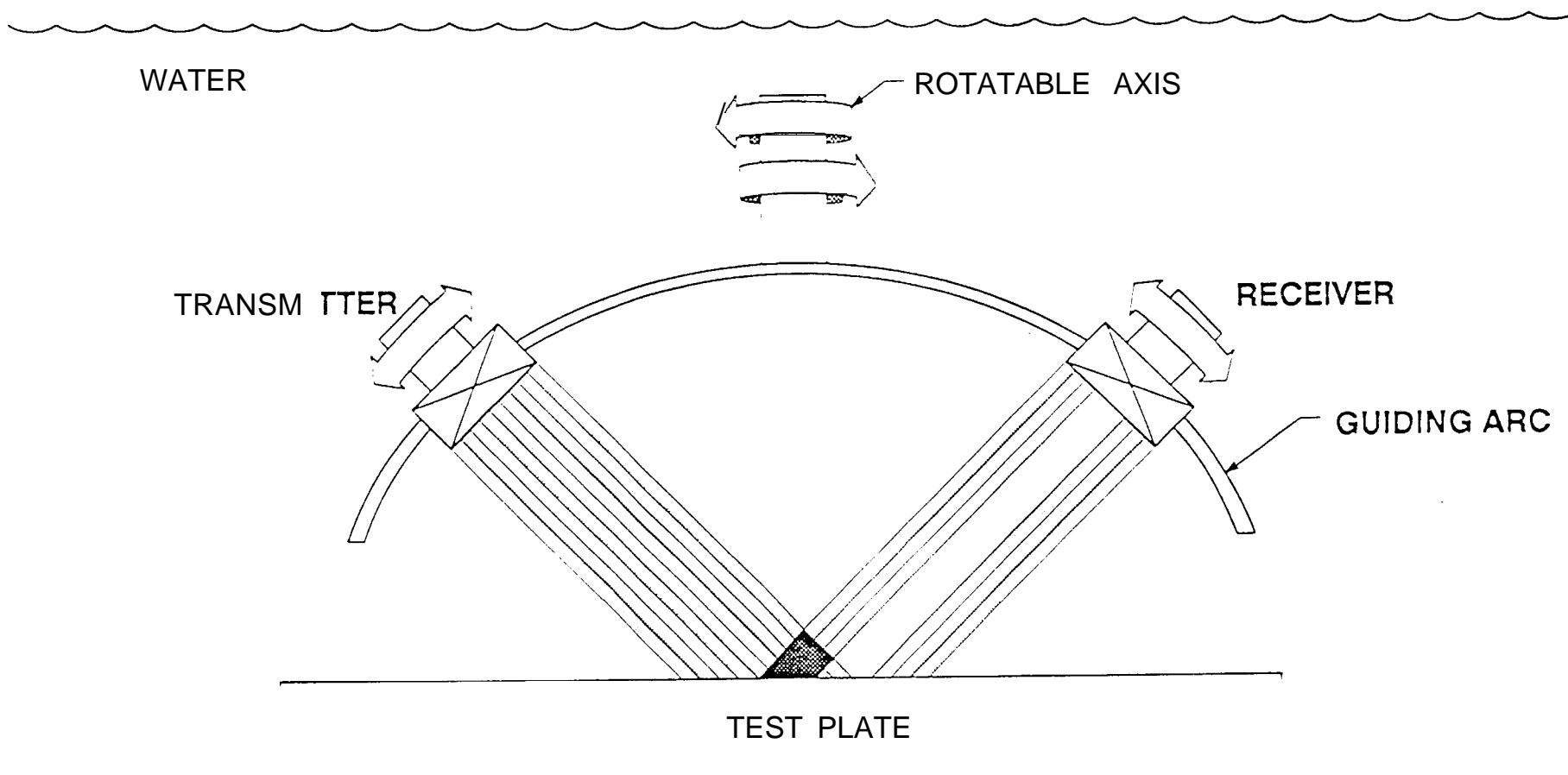
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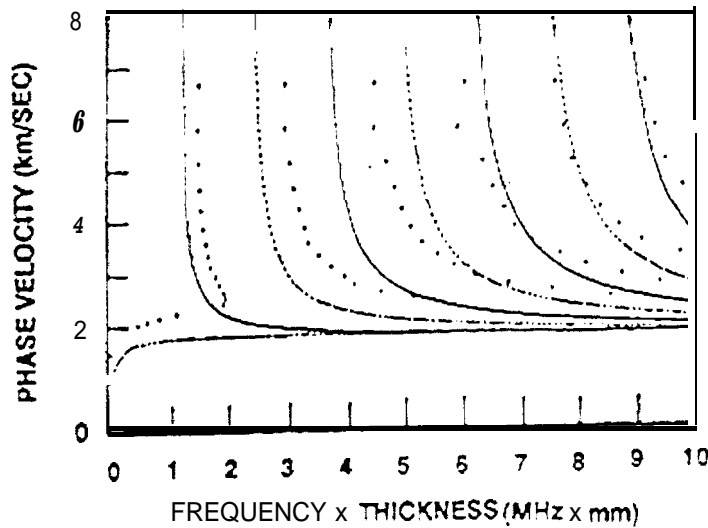
Quick Test

EXPERIMENTAL SETUP FOR COMBINED PBS AND LLW

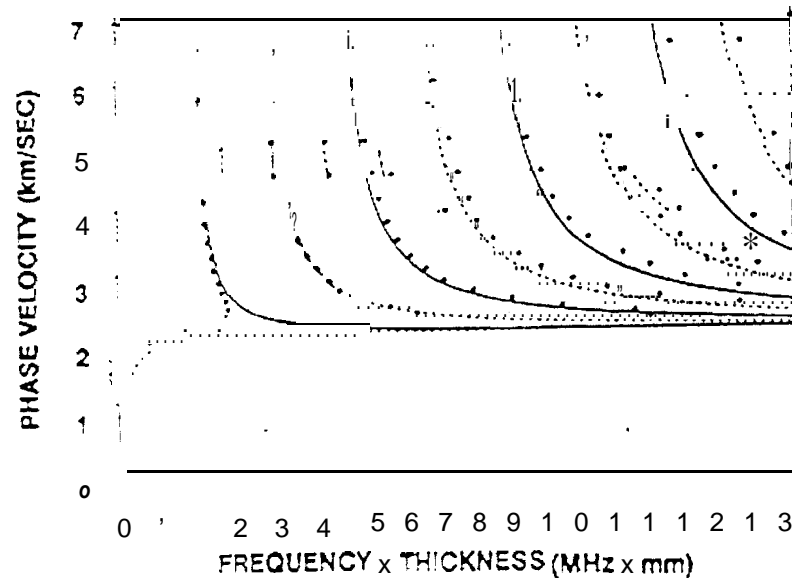


Leaky Lamb Wave (LLW) Dispersion Curve For A Graphite/Epoxy Laminate

a. ASSUMED ELASTIC PROPERTIES



b. INVERTED ELASTIC PROPERTIES



Nonlinear Electroelastic Equations

Stress equations of motion:

$$K_{Lj,L} = \delta_{jM} \rho^o \ddot{u}_M$$

Charge equations:

$$\hat{D}_{L,L} = 0$$

where ρ^o , u_M , K_{Lj} , \hat{D}_L , and δ_{jM} denote the reference mass density, the mechanical displacement, the total Piola-Kirchhoff stress tensor, the reference electric displacement vector, and the translation operator (where capital-reference state, lower case - present state.)

The parameters can be defined as

$$\begin{aligned}\hat{D}_L &= J X_{L,i} D_i, \quad y_i = y_i(X_L, t) \\ J &= \det y_{i,L}, \quad y_i = \delta_{iM}(X_M + u_M) \\ K_{Lj} &= H_{Lj} + M_{Lj}, \quad \hat{D}_L = \epsilon_0 \hat{E}_L + \hat{P}_L\end{aligned}$$

where ϵ_0 is the electric permittivity of free space, and

$$H_{Lj} = \rho_0 v_{j,M} \frac{\partial \chi}{\partial E_{LM}}$$

$$P_L = -\rho^0 \frac{\partial \chi}{\partial W_L}$$

$$M_{Lj} = J X_{Li} T_{ij}^{ES}, E_L = J X_L E_i$$

$$T_{ij}^{ES} = \epsilon_0 (E_i E_j - \frac{1}{2} (y_{iL} y_{iM} - \delta_{LM})), W_L = y_{iL} E_i$$

where T_{ij}^{ES} is the symmetric Maxwell electrostatic stress tensor, E_{LM} is the material strain tensor, and W_L is the rotationally invariant electric variable. The Gibbs free energy function $\chi = \chi(E_{KL}, W_L)$ is defined through

$$\begin{aligned} \rho_0 \chi = & \frac{1}{2} c_{ABCD} E_{AB} E_{CD} - e_{ABC} W_A E_{BC} \\ & - \frac{1}{2} \xi_{AB} W_A W_B - \frac{1}{2} b_{ABCD} W_A W_B E_{CD} \\ & - \frac{1}{6} \xi_{ABC} W_A W_B W_C \\ & - \frac{1}{6} d_{ABCDE} W_A W_B W_C E_{DE} \\ & - \frac{1}{24} \xi_{ABCD} W_A W_B W_C W_D \end{aligned}$$

$$W_L = -\phi_{,L}$$

where c, e, ξ, b are elastic, piezoelectric permeability and electrostrictive constants, respectively.

Equations for Small Strain Elements

Constitutive Equation:

$$\begin{aligned} T_p &= c_{pq} S_q - e_{kp} E_k - \frac{1}{2} \hat{b}_{klp} E_k E_l \\ D_l &= e_{lq} S_q + \epsilon_{lk} E_k + \frac{1}{2} \chi_{kjl} E_k E_j \end{aligned}$$

where the contract notation is used, i.e. 11.22,33,23,31,12 = 1,2,3,4,5,6, respectively. or

$$S_q = s_{qp} T_p + d_{kq} E_k + \frac{1}{2} \beta_{jkq} E_j E_k$$

with

$$\begin{aligned} d_{kp} &= e_{kp} s_{pq}, \beta_{klq} = \hat{b}_{klq} s_{pq} \\ s_{rp} &= c_{rp}^{-1} \end{aligned}$$

where S_q is the contracted strain tensor, and d_{kq} and β_{kl} are the piezoelectric constants and nonlinear electrostrictive constants, respectively.

Equations for Thin Elements (plane stress)

$$T_{3m} = 0, D_{3,3} = 0, D_3 = D_3(X_1, X_2, t)$$

$$E_1 = E_2 = 0, E_3 = -V/t.$$

where V is the driving voltage, and t is thickness of the element. Then the constitutive equations reduced to

$$S_q = s_{qv}T_v + d_{3q} + \frac{1}{2}\beta_{33q}E_3^2$$

$$D_l = d_{lv}T_v + e_{l3}^t E_3 + \frac{1}{2}\chi_{33l}E_3^2$$

where $v=1, 2, 6$, or

$$S_1 = s_{11}T_1 + S_{12}T_2 + d_{31}E_3 + \frac{1}{2}\beta_{31}E_3^2$$

$$S_2 = s_{12}T_1 + S_{11}T_2 + d_{31}E_3 + \frac{1}{2}\beta_{31}E_3^2$$

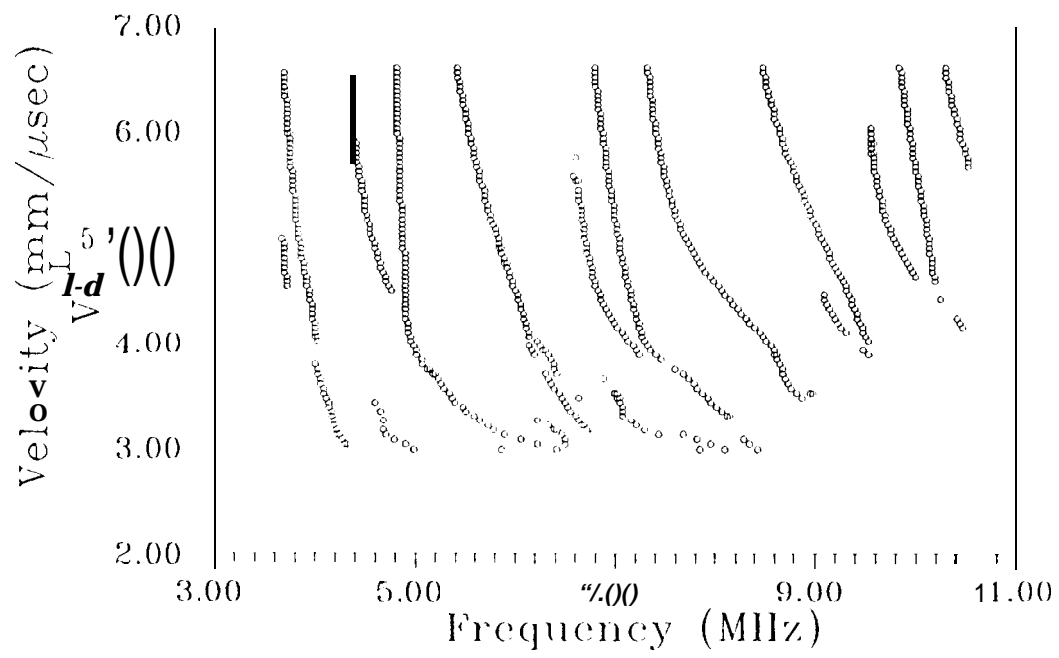
$$S_3 = s_{13}T_1 + S_{13}T_2 + d_{33}E_3 + \frac{1}{2}\beta_{33}E_3^2$$

$$S_6 = s_{66}T_6, S_4 = S_5 = D_1 = D_2 = 0$$

$$D_3 = d_{31}(T_1 + T_2) + \epsilon_{33}^T E_3 + \frac{1}{2}\chi_{33}^T E_3^2$$

Stress-Free Thin Element Subject to Large Electric Fields

$$S_1 = S_2 = d_{31}E_3 + \frac{1}{2}\beta_{31}E_3^2$$



SUMMARY

- Leaky lamb wave measurements were used to determine the thickness and the elastic properties of electroactive thin film polymers.
- Results show that LLW dispersion data can be used to measure thin-films thickness at high resolution using significantly lower frequencies than possible with pulse-echo.
- These measures also allow simultaneous determination of the elastic constants.