

sion of $\sim 4^\circ$. However, Dondi & Ghisellini (1995) argue that the radio and gamma-ray beams must be collimated to the same extent, and they suggest that the EGRET sources are detected because they are currently in a high state of gamma-ray emission, thus leaving those sources currently in a low state undetected.

Indicators of relativistic beaming are therefore relevant for models of the gamma-ray emission. For example, the models proposed by Salamon & Stecker (1994) and Dondi & Ghisellini (1995) have quite different implications for the distribution of line of sight angles of the relativistic jet in EGRET sources. Salamon & Stecker predict that all gamma-ray sources have their jets aligned within 4° to the line of sight, whereas Dondi & Ghisellini imply a broader distribution of angles to the line of sight.

Such predictions of the beaming characteristics of EGRET sources may be tested by comparing the VLBI properties of radio sources, such as superluminal motions and radio core brightness temperatures, over the full range of gamma-ray activity (nondetections, weak detections, and strong detections) and examining any similarities or differences between these populations.

In this paper we describe the first high-resolution VLBI observations of three southern EGRET radio sources, PKS 0208-512, PKS 0521-365, and PKS 0537-441, all high-confidence identifications in the second EGRET cycle. In § 2, our observations, reductions, and analysis methods are outlined. The results for the individual sources are presented in § 3, where we derive estimates of the radio core brightness temperatures. We conclude with a discussion of the existing VLBI observations of EGRET sources in § 4.

2. OBSERVATIONS AND DATA REDUCTION

Our VLBI observations, summarized in Table 1, were all made using the Mark-II recording system (Clark 1973) and the SHVEE (Southern Hemisphere VLBI Experiment) array of telescopes. The array, its properties, and operational modes have been described elsewhere (Preston et al. 1989; Jauncey et al. 1994).

All data were recorded with left circular polarization (right-hand convention) and upper sideband with a bandwidth of 1.8 MHz, and subsequently they were correlated at the Block II Caltech/JPL processor in Pasadena, California. The data were fringe-fitted in NRAO's Astronomical Image Processing System (AIPS) before being passed to the Caltech reduction software package for coherent averaging in time. An initial calibration, available from measured

system temperatures and antenna sensitivities, was applied to calibrator sources of known total flux density observed in the same experiment. A single time-independent scaling factor was determined for each antenna by amplitude self-calibration on the calibrator sources and applied to the amplitudes of the program sources to improve the initial calibration. The error in the flux density scales was then estimated by imaging the program source data and noting the telescope-based corrections made during the self-calibration stages. In all cases, these corrections were less than 10% and often less than 5%. The data were imaged using DM3MAP (Shepherd, Pearson, & Taylor 1994), part of the Caltech package (Pearson 1991) with 512×512 maps and pixel sizes ranging between 0.1 and 0.5 mas. Brightness temperatures were found for components which were deconvolved from the synthesized beam using the JMFIT task in AIPS. JMFIT returns only an upper limit on the size of completely unresolved components but returns an estimated size, with upper and lower limits, for components that are resolved. In the case of unresolved core components, we were able to derive lower limits to the brightness temperature from the limit on the size and the total flux of the component. For the cores that were resolved, we could derive an estimate of the brightness temperature from the estimated size and the total flux.

For PKS 0521-365, high-resolution multiepoch and multifrequency data are available, and so, in addition to the above procedures, the following analyses were performed. Limits on the jet-to-counterjet sulfate brightness ratios were made directly from the images. We estimated spectral indices using the total flux in each component measured from the images. The multiepoch data were used to estimate the apparent speed in the jet of PKS 0521-365. At this end, we used the Caltech task MODELFIT to fit simple Gaussian components to the amplitudes and closure phases, either two or three components at each epoch, keeping the core component as our arbitrary phase center.

3. THE INDIVIDUAL SOURCES

3.1. PKS 0208-512

PKS 0208-512 was detected with a strong statistical significance ($> 5\sigma$) by EGRET (von Montigny et al. 1995). The high-confidence identification is with an $m_r = 17.5$ high-polarization quasar at a redshift of 1.003 (Savage 1975; Peterson et al. 1976; Impey & Tapia 1988). PKS 0208-512 has the hardest gamma-ray spectrum of all EGRET sources, with a differential spectral index of 1.689 ± 0.06

TABLE 1
OBSERVATION LOG

Source	Epoch	Frequency (GHz)	Participating Telescopes*
PKS 0208-512	1992 Nov 28	4.851 GHz	Pk, Hb, Cg, M1, Pr27, Ht
PKS 0521-365	1992 Nov 23	4.851 GHz	Pk, Jib, Cg, M1, J17, Ht, Sh
	1993 Feb 15	4.851 GHz	Pk, Hb, Cg, M1, Pr27
	1993 May 14	4.851 GHz	Pk, Hb, M1, Pr27
	1993 Oct 21	8.418 GHz	Pk, Hb, Cg, M1, Pr15
PKS 0537-441	1992 Nov 24	4.851 GHz	Pk, Hb, Cg, M1, Pr27, Ht, Sh
	1994 Feb 25	8.418 GHz	Ds45, Pk, Hb, Cg, M1

* Ds45 = Tidbinbilla (Deep Space Network, 34 m), Pk = Parkes (Australia Telescope National Facility [ATNF], 64 m), Hb = Hobart (University of Tasmania, 26 m), Cg = Culgoora (ATNF, 22 m), M1 = Mopra (ATNF, 22 m), Pr15 = Perth (European Space Agency, 15 m), Pr27 = Perth (Telstra, 27 m), Ht = Hartbeesthoek Radio Astronomy Observatory (26 m), Sh = Shanghai Astronomical Observatory (25 m)

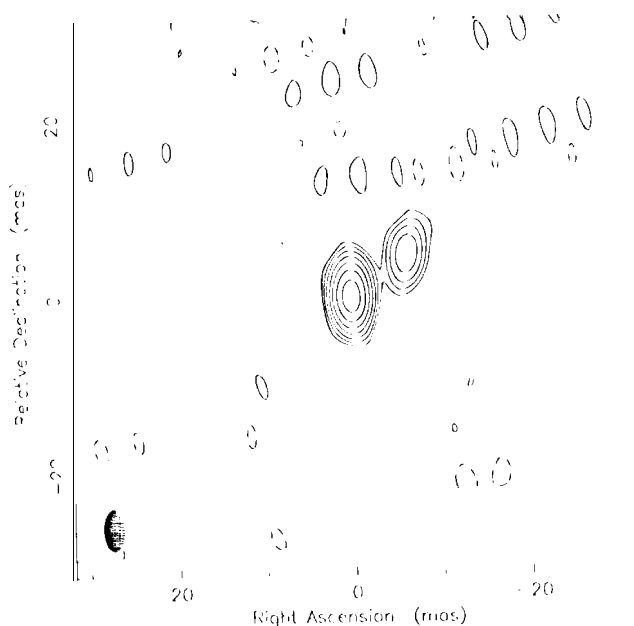


FIG. 2. VLBI image of PKS 0521-365 at 4.851 GHz from 1992 November 23. Map peak is $1.27 \text{ Jy beam}^{-1}$ with contours of 1, 1.2, 4, 8, 16, 32, and 64% of peak and a beam FWHM of 4.68×2.32 (mas) at a position angle of 1.45° .

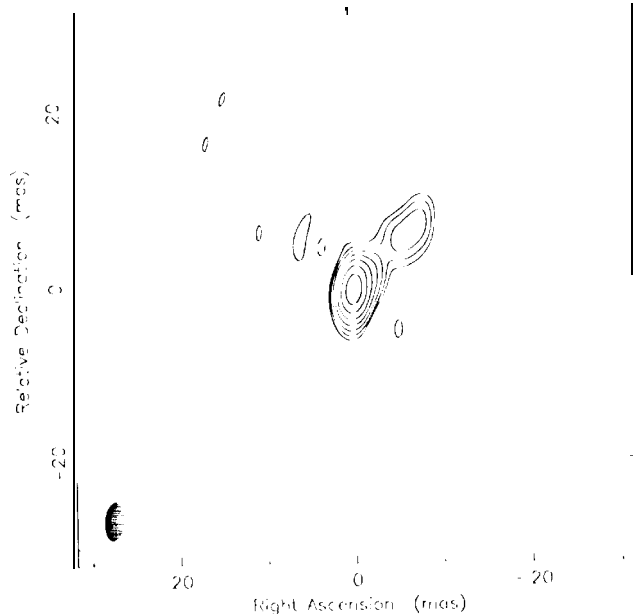


FIG. 4. VLBI image of PKS 0521-365 at 4.851 GHz from 1993 May 14. Map peak is $1.95 \text{ Jy beam}^{-1}$ with contours of 1, 1.2, 4, 8, 16, 32, and 64% of peak and a beam FWHM of 4.26×2.13 (mas) at a position angle of -6.29° .

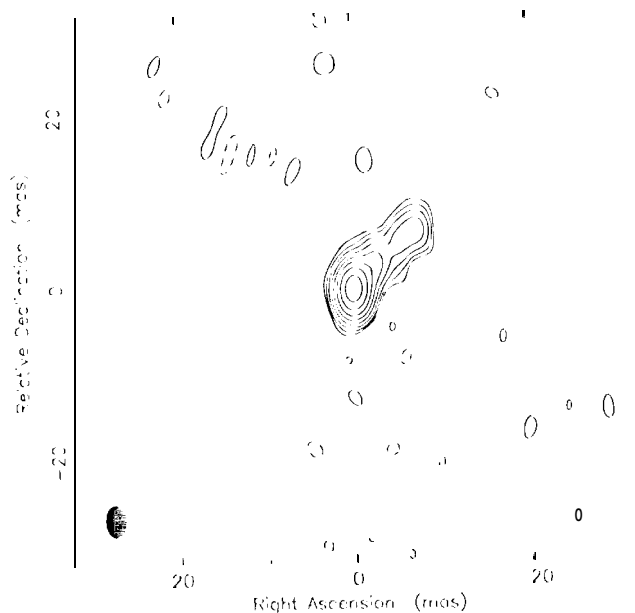


FIG. 3. VLBI image of PKS 0521-365 at 4.851 GHz from 1993 February 15. Map peak is $1.77 \text{ Jy beam}^{-1}$ with contours of 0.5, 0.5, 1, 2, 4, 8, 16, 32, and 64% of peak and a beam FWHM of 3.74×2.27 (mas) at a position angle of 0.58° .

proposed by Falomo et al. (1995) for PKS 0521-365 but also with a model which has a highly relativistic flow closely aligned to our line of sight. However, the brightness temperature, which is well below the nominal inverse Compton limit of 10^{12} K, does not favor PKS 0521-365 as a highly beamed radio source. Observations at later epochs will allow us to constrain further any apparent motion in the jet of PKS 0521-365, giving a more complete indication of the importance of relativistic beaming.

TABLE 2
LIMITS ON JET-TO-COUNTERJET SURFACE BRIGHTNESS
RATIO (R) AND CORE-JET SEPARATION (IN μ) FOR EACH
OF THE PKS 0521-365 EPOCHS

Epoch	Frequency (GHz)	R	μ (mas)
1992.91.....	4.8	> 9.3	7.8 ± 1.4 1.6
1993.1 ?.....	4.8	> 8.6	8.6 ± 1.6 1.6
1993.3 J.....	4.8	> 5.0	9.4 ± 1.4 2.0
1993.80.....	8.4	> 2.9	10.3 ± 1.8 1.8

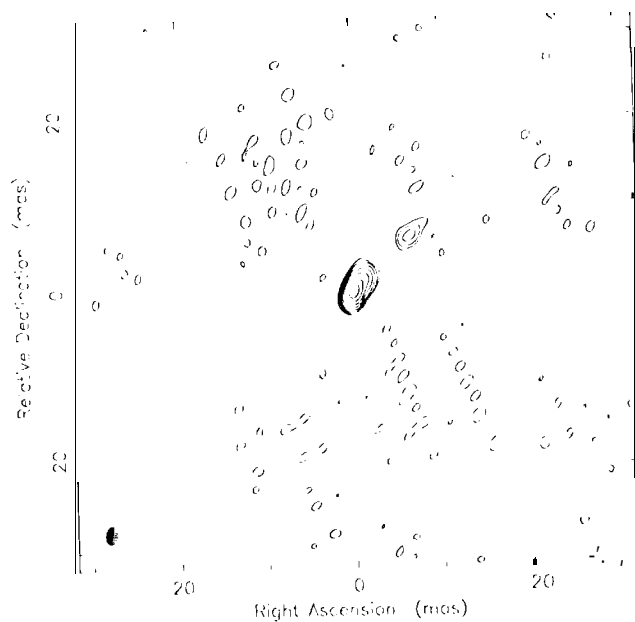


FIG. 5. VLBI image of PKS 0521-365 at 8.418 GHz from 1993 October 21. Map peak is $1.56 \text{ Jy beam}^{-1}$ with contours of 1, 1.2, 4, 8, 16, 32, and 64% of peak and a beam FWHM of 2.04×1.34 (mas) at a position angle of -10.1° .

core-jet morphologies. In the case of PKS 0208-512 and PKS 0521-365, the position angle of the millisecond-scale structure seen with VLBI aligns with the position angle of the arcsecond-scale radio structure. The millisecond-scale position angle of PKS 0537-441 is somewhat misaligned from its arcsecond-scale position angle.

The brightness temperatures of the cores suggest that PKS 0208-512 and 0537-441 may be more highly beamed than PKS 0521-365. In particular, PKS 0208-512 has a brightness temperature at the 10^{12} K inverse Compton limit for synchrotron radiation. This difference in brightness temperature correlates with the strong statistical EGRET detection of PKS 0208-512 and 0537-441 ($> 5\sigma$ significance) but only the weak statistical detection of PKS 0521-365 ($4 < \sigma < 5$ significance), even though PKS 0521-365 is much closer to us. It is important to repeat here that we are considering the significance of the EGRET detection based on the photon statistics, given that the identifications can be made, as Thompson et al. (1995) assert, with high confidence.

Unfortunately, on the basis of the brightness temperatures for these three sources we cannot make a strong statement concerning the importance of relativistic beaming for detection by EGRET. Further constraints for these objects will be available as we obtain images from future epochs. Can we, however, examine the existing VLBI data in the literature and find any trends which could link the gamma-ray and VLBI characteristics of the EGRET-identified sources or differentiate the EGRET-identified radio sources from radio sources which have not been identified by EGRET?

First, nine of the 11 identified EGRET sources listed by Thompson et al. (1995) which have been well studied with VLBI are known or possible superluminal radio sources (Vermeulen & Cohen 1994; Barthel et al. 1995; this work). The mean apparent speed in the VLBI jets of these 11 EGRET-identified sources is $5.6h^{-1}c$ with a standard deviation of $7.3h^{-1}c$.¹

Second, while there is substantial evidence that some strong EGRET-identified radio sources are highly beamed,

¹ This estimate includes the highest reliability VLBI observations listed by Vermeulen & Cohen (1994) for nine sources, the observations of Barthel et al. (1995) for 1633+382, and the data for PKS 0521-365 appearing in Table 2. Where multiple observations of similar reliability were available for a single source from Vermeulen & Cohen (1994), they were used to form an average value for the source.

some others do not appear to be. Some of the strong EGRET sources, for example 1226-1023 (e.g., Unwin et al. 1985), are clearly highly superluminal, whereas others such as 1222-1216 (Hoommeyer et al. 1992) are perhaps consistent with subluminal speeds. There are also some highly beamed radio sources that have not been identified by EGRET, such as 1641-1399 (e.g., Biretta, Moore, & Cohen 1986). The mean apparent speed for the core selected quasars and BL Lac objects which are listed by Vermeulen & Cohen (1994) and have not been identified by EGRET is $3.9h^{-1}c$, with a standard deviation of $3.0h^{-1}c$ (based on 24 objects and using the same criteria as for the EGRET-identified sources).

Unfortunately, these data do not allow us to make a firm conclusion about possible differences or similarities in VLBI properties between the radio sources that EGRET has identified and the many others that EGRET has not identified. The mean values for the apparent VLBI speeds are easily consistent at the 1σ level, but the standard deviations on the distributions are large. We feel that, given the overall reliability of the data and the small samples, any definite conclusion is premature.

Thus, more VLBI observations, over the full range of gamma-ray strength (non-detection, weak detection, and strong detection), will be required before the question posed above can be properly answered. The need for Southern Hemisphere observations of the high-confidence EGRET identifications is especially apparent since 13 of 40 lie south of $\delta = 0^\circ$. Our continuing VLBI observations are aimed at addressing this need.

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