

# Recent results with coupled opto-electronic oscillator

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## ABSTRACT

We present experimental results of coupled opto-electronic oscillators (COEO) constructed with a semiconductor optical amplifier based ring laser, a semiconductor Fabry-Perot laser, and a semiconductor colliding pulse mode-locked laser. Each COEO can simultaneously generate short optical pulses and spectrally pure RF signals. With these devices, we obtained optical pulses as short as 6 picoseconds and RF signals as high in frequency as 18 GHz with a spectral purity comparable with a HP8561B synthesizer. These experiments demonstrate that COEOs are promising compact sources for generating low jitter optical pulses and low phase noise RF/millimeter wave signals.

Key words: Opto-electronic oscillator, optical pulse, RF photonics, signal generation, mode-locked lasers, microwave generation, phase noise, jitter.

## 1. INTRODUCTION

Opto-electronic oscillators (OEO)<sup>1,2</sup> are a new class of oscillators for generating high spectral purity, high frequency, and high stability RF signals and optical subcarriers. We have demonstrated an OEO operating at 10 GHz with a phase noise of -140 dBc/Hz @ 10 kHz from the carrier.<sup>3</sup> This is the lowest phase noise signal from a free running oscillator operating at room temperature.

In previously reported OEOs, the optical oscillation of the pump laser is isolated from the electronic oscillation. In a recent paper,<sup>4</sup> we demonstrated a Coupled Opto-Electronic Oscillator (COEO) in which the laser oscillation is directly coupled with the electronic oscillation. The coupling of the microwave and optical oscillations causes the laser to modelock, generating stable optical pulses and microwave signals simultaneously. Because of its unique features, COEOs will find wide applications in RF communication systems, fiber optic communication systems, and photonic analog-to-digital conversion systems.

In this paper, we present three new configurations of COEO for simultaneously generating low jitter picosecond optical pulses and high spectral purity microwave signals beyond 10 GHz. We will discuss, in particular, two configurations involving integration of the laser and modulator to make compact and cost effective COEOs.

## 2. RING LASER BASED COEO

In the previously reported COEO,<sup>4</sup> a semiconductor optical amplifier (SOA) is used in a ring configuration to form the optical oscillation loop and the RF oscillation loop directly feeds back to the SOA to modulate its gain. Due to the slow response of the SOA, the RF oscillation frequency is limited to below 1 GHz. To increase the RF oscillation frequency, here we use a Mach-Zehnder modulator in the laser oscillation loop to modulate the loop gain.

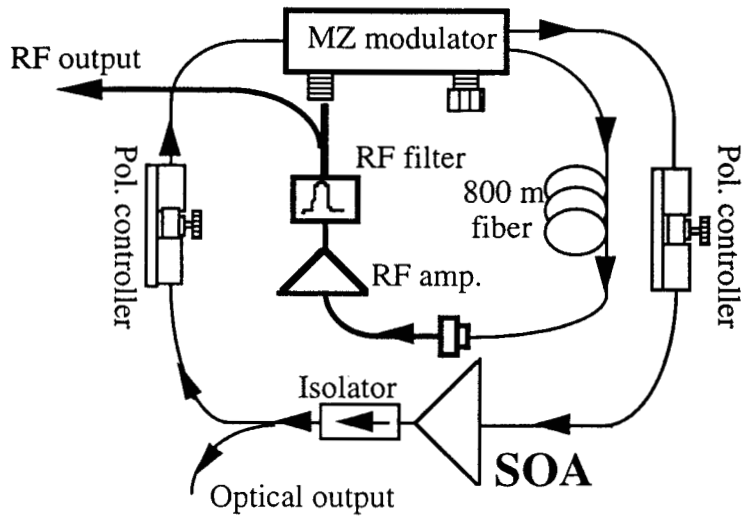


Fig. 1 The experimental setup of a COEO constructed with a SOA and an external modulator.

The experimental setup of the COEO is shown in Fig. 1. The output of a semiconductor optical amplifier (SOA) is connected to a Mach-Zehnder modulator with a 10 GHz bandwidth. One of the outputs from the modulator is fed back to the SOA via a polarization controller to form a ring laser. The other output port of the modulator is delayed by an 800 meter optical fiber, detected by a photodetector, amplified by an RF amplifier, filtered by an RF bandpass filter centered at 10 GHz, and finally is coupled to the RF modulation port of the modulator to form an opto-electronic feedback loop. Just like an OEO, when the gain of the feedback loop is larger than one, an opto-electronic (O/E) oscillation will start. As described in reference 4, the interaction between the optical oscillation modes of the ring laser and the O/E oscillation modes will force the laser to mode-lock. The mode-locked laser will in turn reinforce the O/E oscillation.

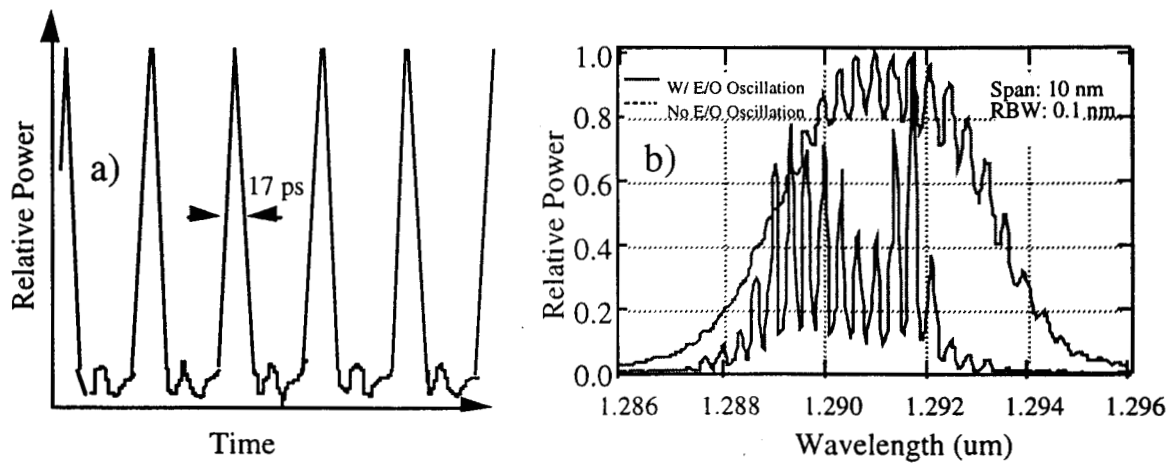


Fig. 2 a) Time domain measurement of the optical pulses. b) Optical spectrum of the ring laser with and without opto-electronic oscillation.

The experimental arrangement is similar to a regeneratively mode-locked laser;<sup>5</sup> however in our case the O/E oscillation modes, which were not considered in regenerative mode-locking, play a critical role. In particular, the long fiber delay for the opto-electronic loop stores the phase information of both the opto-electrical oscillation and the optical oscillation. The feedback of the stored phase information is the key to high spectral purity oscillations. Similar to a conventional OEO, the phase noise of the O/E oscillation, which directly translates to the jitter of the optical pulses, is expected to be inversely proportional to the time delay squared. Longer O/E loop delay reduces the phase noise of the generated microwave and lowers the jitter of the optical pulses.

The pulses generated by the COEO were measured with a New Focus 40 GHz detector and a Tek CSA803 communication signal analyzer, and the result is shown in Fig. 2a. The measured pulse width is 17 ps, limited by the rise time of the sampling head (SD-26). However, our preliminary autocorrelation measurement indicated a pulse width of 15 ps. On the other hand, the optical spectrum of the pulses, as shown in Fig. 2b, has a bandwidth of 4 nm, implying that the pulses were not transform limited.

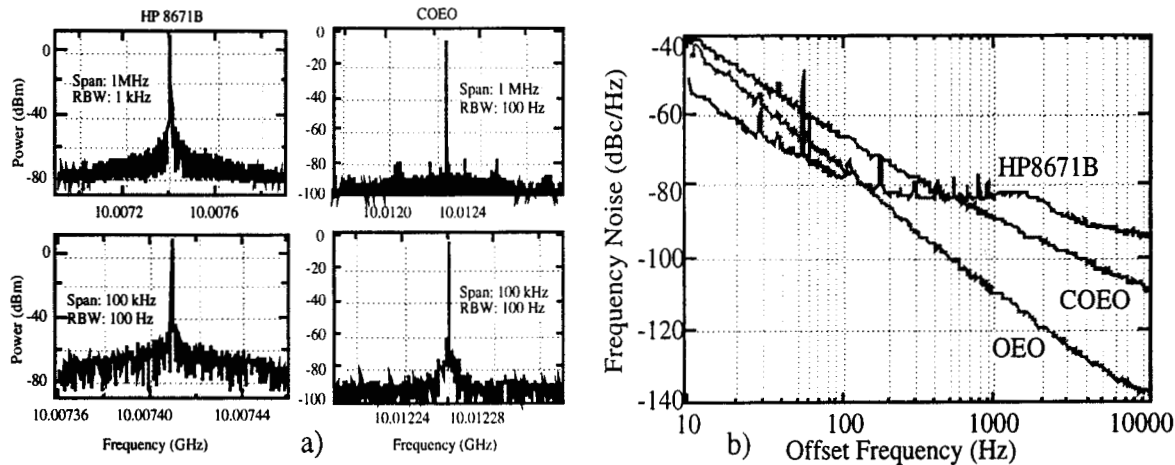


Fig. 2 a) Measured RF spectra of the COEO output at different spectrum analyzer settings and the comparison with a signal from a HP8671B synthesizer. b) Single sideband phase noise comparison of the COEO, an OEO and a HP 8671B synthesizer at 10 GHz.

The measured RF spectrum are compared with a high performance synthesizer (HP8671B), as shown in Fig. 3a. Clearly, for the spectrum analyzer settings, the spectral purity of the COEO is better than that of HP8671B. We also measured the phase noise of the COEO and the result is shown in Fig. 3b. For comparison, the phase noises of HP8671B synthesizer and a conventional OEO with 2 km loop length are also shown in Fig. 3b. It is evident that at 10 kHz from the carrier, the phase noise of the COEO is about 10 dB better than that of HP8671B, however, substantially larger than that of the conventional OEO. We expect to further lower the phase noise of the COEO by increasing its loop delay and employing other noise reduction techniques.

### 3. FABRY-PEROT LASER-BASED COEO

A COEO can also be constructed with Fabry-Perot lasers for reduced size and cost, as shown in Fig. 4. To increase optical power and modulation speed, and reduce the chirp effect, an active gain medium is integrated with an electro-absorption

modulator inside the laser cavity. Because electroabsorption modulators with a bandwidth greater than 60 GHz have been demonstrated, such an approach is promising for achieving millimeter wave oscillations.

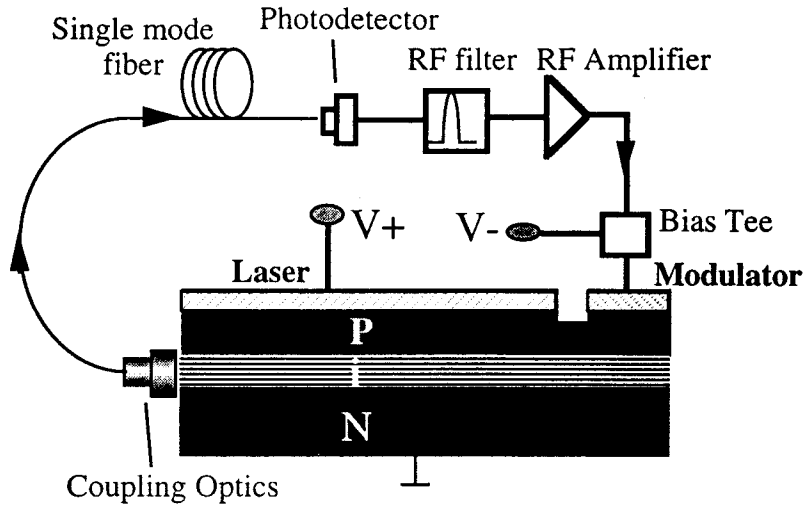


Fig. 4 Illustration of a COEO constructed with a Fabry-Perot laser and an integrated electro-absorption modulator inside the laser cavity.

The integrated laser/modulator uses the identical active layer approach. The active layer of the laser and modulator is made of InGaAsP/InP multiple quantum wells with a graded index InGaAsP cladding. The total length of the integrated laser/modulator is about 3 mm, which gives a Fabry-Perot mode spacing of 15 GHz. The laser and modulator are electrically isolated by an etched groove in between them. The modulator has a length of less than 100  $\mu$ m. The lasing wavelength is around 1350 nm. When no bias is applied to the modulator section, the laser shows a threshold current of 50 mA. The facet output power from the laser is about 10 mW at a driving current of 150 mA, and 20 mW at driving current of 300 mA. A tapered single mode fiber is used for butt coupling and the coupling efficiency is more than 50%. Fig. 5 shows the laser output power (coupled into a single mode fiber) as a function applied voltage across the modulator. The high output power of this laser is excellent for obtaining OE oscillation.

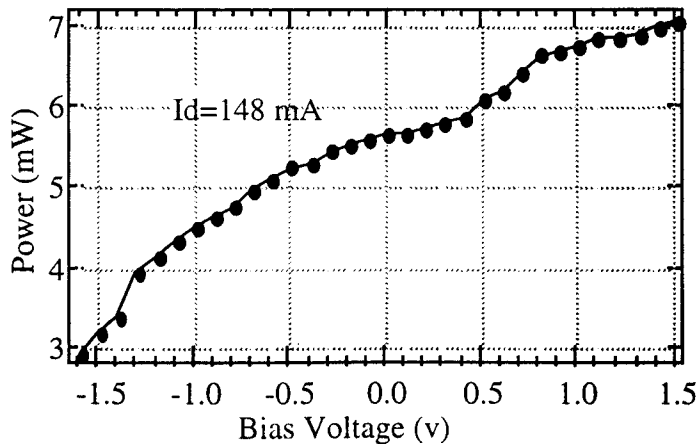


Fig. 5 The laser output power coupled into a single mode fiber as a function of bias voltage across the electro-absorption modulator.

The chip is bonded to a Silicon carrier for RF testing. The RF reflection and modulation response were measured with a HP8607A Lightwave Network Analyzer and are shown in Fig. 6a and 6b. As indicated in Fig. 6a, the modulation response has a sharp peak at 13.6 GHz, due to the resonant enhancement. When closing the OE loop, this resonant enhancement will force the OE oscillation at 13.6 GHz and cause the laser to mode-lock, as explained in the first section. As can be seen, the reflection coefficient of the device is about -15 dB at 10 GHz and -8 dB at 15 GHz.

The mode-beating spectrum of a free running laser was observed with a photodetector directly connected to a RF spectrum analyzer and is shown in Fig. 7a. The strong mode beating at 13.6 GHz again indicates that when closing the OE loop, the COEO will be forced to oscillate at 13.6 GHz and cause the laser to mode-lock. The RF spectrum also indicates that the laser is free of self-pulsation under the operating condition.

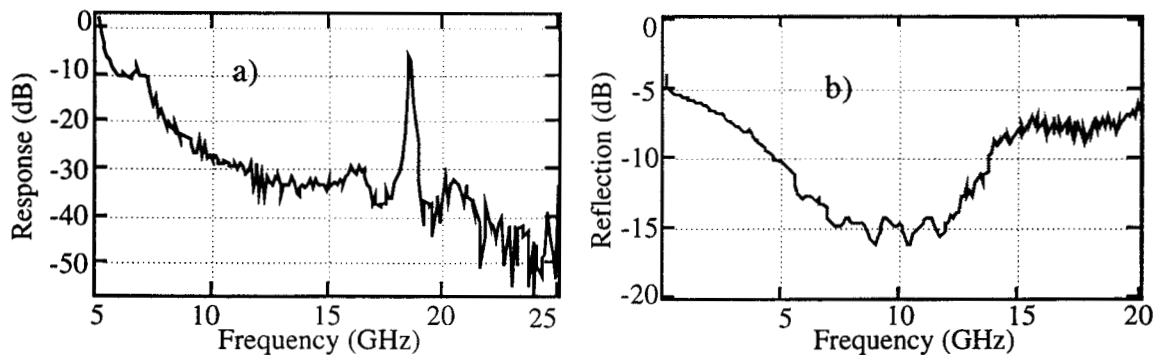


Fig. 6 a) RF response of the integrated laser/modulator. b) RF reflection coefficient of the integrated laser/modulator.

The phase noise of the mode-beating signal was significantly reduced when a RF signal of 5 dBm at 13.6 GHz was applied to the E/A modulator, as shown in Fig. 7b. This indicates that the laser can easily be mode locked. Unfortunately, due to lacking a key RF component at the time of experiment, we were not able to close the O/E loop and demonstrate COEO operation using this device.

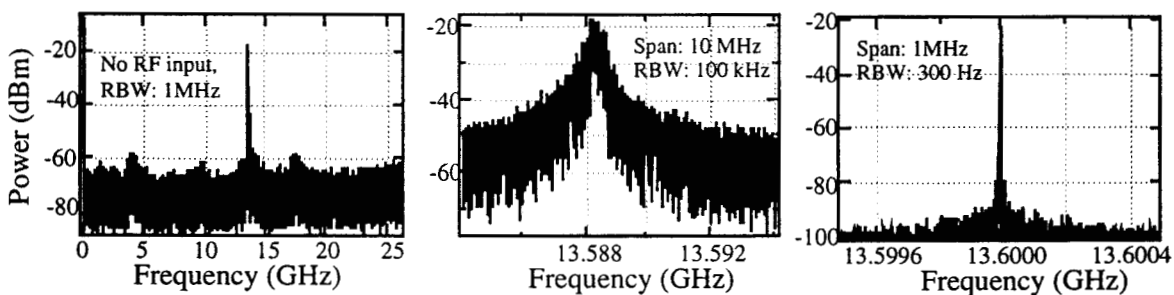


Fig. 7. a) and b) Mode-beating spectrum of a free running integrated laser/modulator for different spectrum analyzer settings. c) The mode beating spectrum of same device with a 5 dBm RF signal at 13.6 GHz applied to the modulator. Note the difference in total span and resolution bandwidth (RBW) of the spectrum analyzer settings.

#### 4. COEO BASED ON COLLIDING PULSE MODE-LOCKED (CPM) LASERS

Colliding pulse mode-locking has been the most effective technique to generate ultrashort optical pulses in passive mode-locked dye lasers. Y. K. Chen and M. C. Wu<sup>6</sup> successfully demonstrated monolithic integration of a CPM laser on a InP substrate and obtained subpicosecond pulses at repetition frequencies up to 350 GHz. We demonstrate here that by

incorporating an electro-optic oscillation loop with a CPM laser, one can greatly reduce the phase noise and frequency jitter of the laser pulses. The effect is similar to that of injection locking the laser with an external RF source, however the external RF source is not needed here and the generated signal quality is not limited by the external signal source. Therefore, such a CPM-based COEO can be used as a stand-alone compact source for both mm-wave signals and subpicosecond optical pulses.

The monolithic CPM laser used in this experiment has been cleaved from a multi-wavelength laser array designed for soliton communications.<sup>7,8</sup> The monolithic CPM laser is fabricated using two MOCVD growths. The first growth is a separate confinement heterostructure with 4 compressively strained ( $\epsilon=1\%$ ) quantum wells at  $1.55\ \mu\text{m}$  and confined on either side by  $120\ \text{nm}$  of InGaAsP ( $\lambda=1.2\ \mu\text{m}$ ). After the diffraction gratings were written by direct write electron beam lithography and etched into the SCH region, the upper cladding and contact layers were grown. The lasers were fabricated into a  $3.5\ \mu\text{m}$  wide ridge laser structure with a continuous active region. As seen in Fig. 8, the symmetric  $4.6\ \text{mm}$  long ridge was divided into 5 sections with 3 contacts: the two end sections are  $75\ \mu\text{m}$  each and contain the gratings; the two gain sections are  $2180\ \mu\text{m}$ ; and a center saturable absorber section of  $50\ \mu\text{m}$ . The sections are separated by four  $10\ \mu\text{m}$  wide gaps, and electrical isolation of  $\sim 1\text{k}\Omega$  is achieved between the sections by etching the InGaAs contact layer. The fabrication of a microwave ground-signal-ground (GSG) contact for the saturable absorber allows for high frequency probing. A continuous gain region has been employed for ease of fabrication and the elimination of multiple reflections within the cavity. The individual devices have been fully packaged in a single sided Butterfly package, including fiber coupling, K-connector for RF input, TEC/thermistor, and DC leads for the gain and grating contacts.

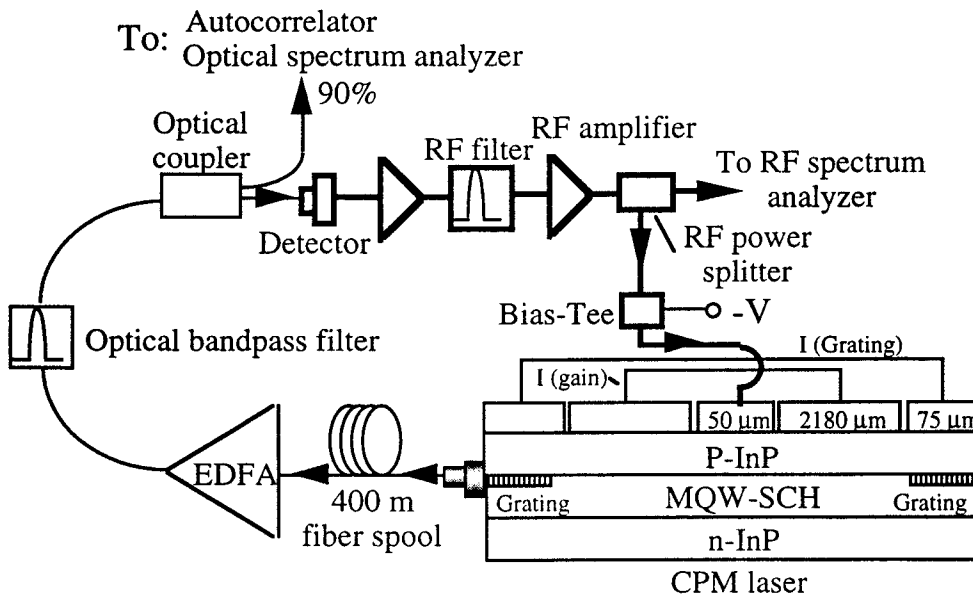


Fig. 8 The cross-section illustration of a colliding pulse mode-locked laser and a COEO constructed with the laser.

The threshold current for uniformly pumped devices without gratings is approximately  $120\ \text{mA}$ . Devices with gratings show thresholds of  $135 - 155\ \text{mA}$  depending on the lasing wavelength with respect to the gain peak. When a reverse bias is applied to the saturable absorber, the devices exhibit passive mode-locking for a range of gain currents and saturable absorber voltages. Typical operating currents are in the range  $165\text{-}210\ \text{mA}$  and typical saturable absorber voltages are  $-0.5$  to  $-2.0$

volts. The devices mode-lock best near threshold (as seen in [6]), and have facet powers of  $\sim 1\text{mW}$  at the optimal operating points. Outside the mode-locking regime of operation, the RF spectra shows a distinct peak at both the cavity fundamental of 9.03 GHz and at the mode-locking frequency of 18.06 GHz. Under the proper bias conditions, the device mode-locks at 18.06 GHz with removal of the CW component, and the 9.03 GHz peak is strongly ( $> 30\text{ dB}$ ) suppressed. However, just like other passively mode-locked lasers, the pulse jitter and phase noise are extremely high due to the high spontaneous emission noise of the laser, the complex interaction of the gain-index-carrier density in semiconductors, and insufficient Q of the laser cavity.

One may supply a sinusoidal clock signal at 18.06 GHz to the saturable absorber of the device to reduce the phase noise and jitter. However, this adds cost, size, and power of the device significantly. In addition, the phase noise of the laser will be limited by the external source.

To make a stand alone low noise signal source, we constructed a COEO with the CPM laser, as shown in Fig. 8. The fiber spool has a length of 400 meters, corresponding to a RF Q of  $2.26 \times 10^5$  for a 18 GHz signal.<sup>2</sup> Note that the use of an EDFA in the loop is not mandatory for the operation of the COEO. It is used merely to boost the optical signal so that the optical pulses can be measured with an autocorrelator. We made the COEO operational even without the insertion of the EDFA in the loop.

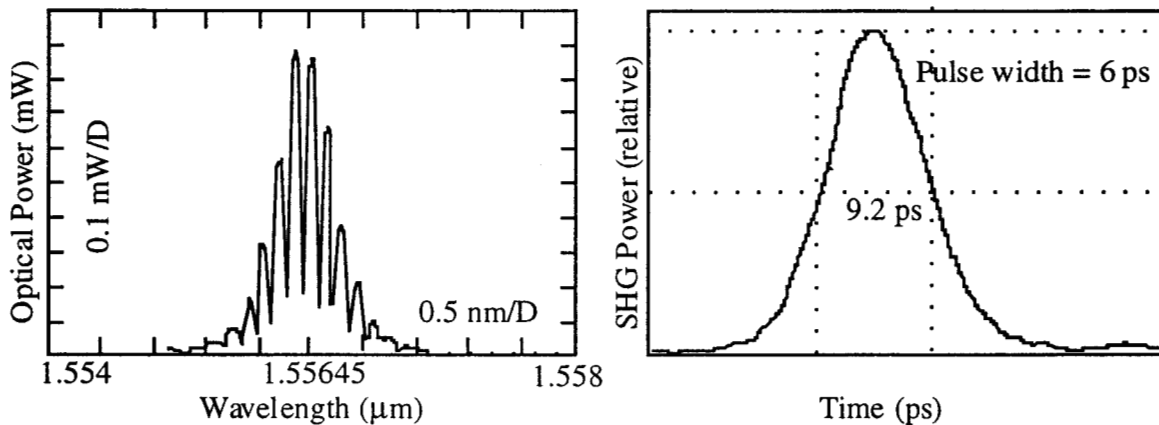


Fig. 9 a) The optical spectrum of a pulse train from a CPM laser based COEO. b) the auto-correlation measurement of the pulse train. Pulse width is 65% of the autocorrelation width if  $\text{sech}^2$  pulse shape is assumed.

After closing the opto-electronic loop, stable mode-locked pulses are immediately present. The spectral and autocorrelation measurements of the optical pulses are shown in Fig. 9a and Fig. 9b respectively. The spectral width  $\Delta\lambda$  is about 0.56 nm and the pulse width  $\Delta\tau$  is 6 ps if  $\text{sech}^2$  pulse shape is assumed.<sup>9</sup> The time and bandwidth product is thus 0.42, slightly above the transform limit of 0.32.

Stable opto-electronic oscillation was also observed with a RF spectrum analyzer at the RF output port of the COEO, as shown in Fig. 10a and 10b. The mode spacing of the O/E oscillation is about 487 kHz, consistent with the O/E loop length. The sidemode suppression is about 21 dB. The single sideband phase noise of the 18 GHz signal can be estimated from the spectrum analyzer measurement to be -104 dBc/Hz at 100 kHz offset and -86 dBc/Hz at 10 kHz offset. The

measured RF spectra of the COEO with or without the EDFA in the loop are about the same. Further phase noise reduction can be achieved by increasing the O/E loop length and by further suppressing the sidemodes using the multi-loop technique.<sup>3</sup>

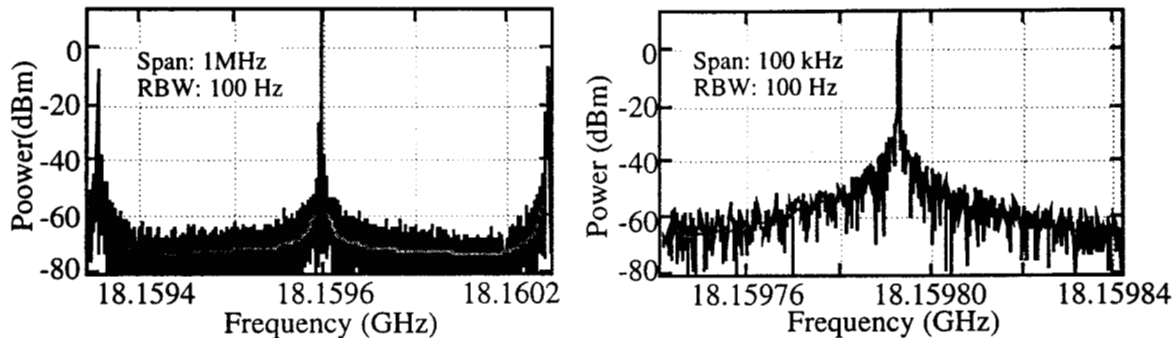


Fig. 10. RF spectrum measurements of a COEO constructed with a CPM laser. Note the different span and resolution bandwidth (RBW) settings of the spectrum analyzer in a) and b).

## 5. SUMMARY

Three new types of coupled opto-electronic oscillators were investigated experimentally. We generated 17 ps optical pulses and a 10 GHz RF signal with low phase noise using a COEO constructed with an SOA ring laser and a Mach-Zehnder modulator. We also demonstrated the generation of 6 ps optical pulses and 18 GHz RF signal with a COEO based on colliding pulse mode-locked (CPM) lasers. Finally, we demonstrated that a high power integrated laser/modulator is promising for making a COEO on a chip, greatly reducing the size, power, and cost of the device. The full demonstration of a COEO based on an integrated laser/modulator is currently under way. We anticipate that the phase noise and jitter of a COEO will reach those of an OEO (phase noise of -140 dBc/Hz @ 10 kHz away from a 10 GHz carrier) in the near future.

## 6. ACKNOWLEDGMENT

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